

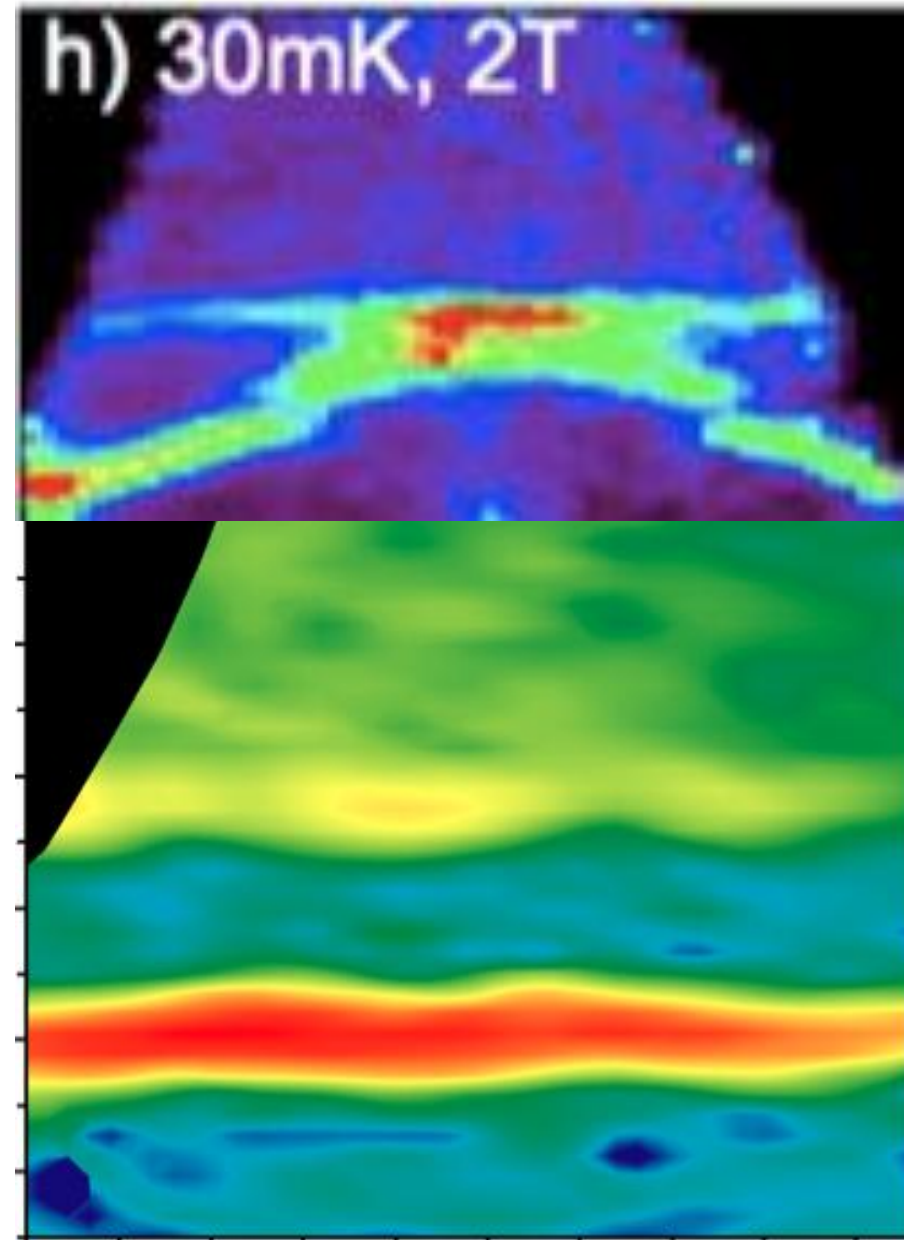
Introduction to Inelastic Neutron Scattering

Bruce D Gaulin
McMaster University



Brockhouse Institute
for **Materials Research**

- *Neutrons: Properties and Cross Sections*
- *Excitations in solids*
- *Triple Axis and Chopper Techniques*
- *Practical concerns*



HAMILTON
ONTARIO

MCMASTER UNIVERSITY

PHYSICS



→

daughter nuclei +
2-3 n + gammas

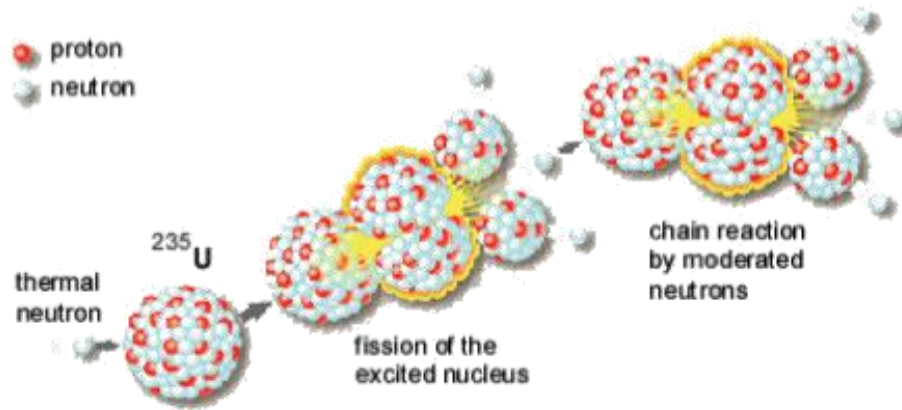
neutrons:

no charge

$s=1/2$

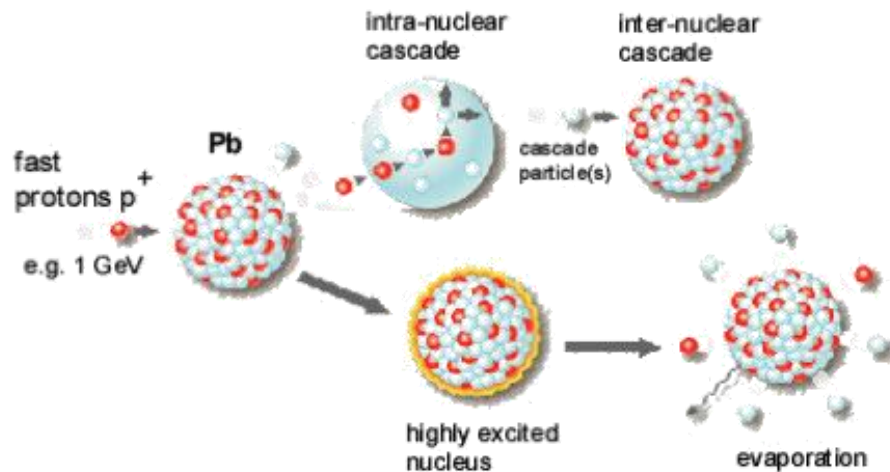
massive: $mc^2 \sim 1 \text{ GeV}$

How do we produce neutrons



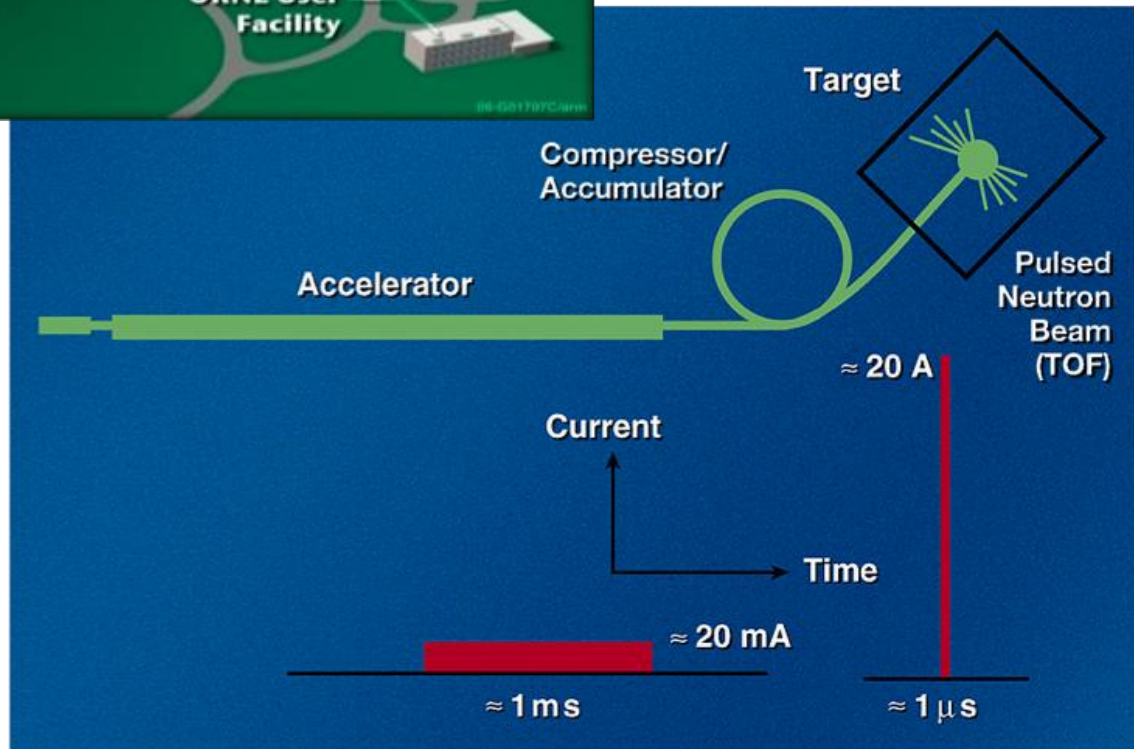
Fission

- chain reaction
- continuous flow
- 1 neutron/fission



Spallation

- no chain reaction
- pulsed operation
- 30 neutrons/proton



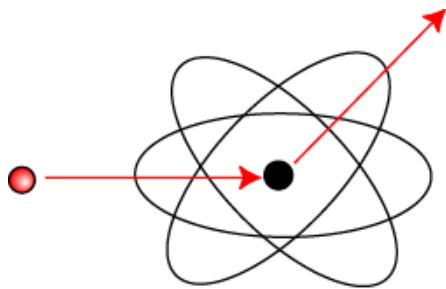
Neutron interactions with matter

- **Properties of the neutron**

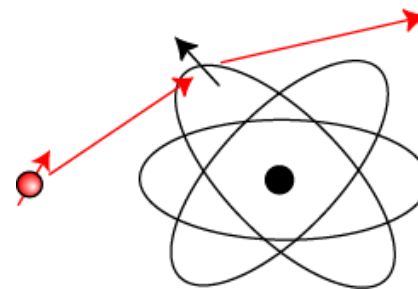
- Mass $m_n = 1.675 \times 10^{-27}$ kg
- Charge 0
- Spin-1/2, magnetic moment $\mu_n = -1.913 \mu_N$

- **Neutrons interact with...**

- Nucleus
- Crystal structure/excitations (eg. Phonons)
- Unpaired electrons via dipole scattering
- Magnetic structure and excitations



Nuclear scattering



Magnetic dipole scattering

Wavelength-energy relations

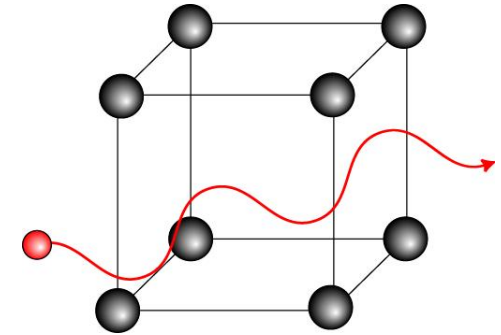
- **Neutron as a wave ...**

- Energy (E), velocity (v), wavenumber (k), wavelength (λ)

$$k = \frac{m_n v}{\hbar} = \frac{2\pi}{\lambda}$$

$$E = \frac{\hbar^2 k^2}{2m_n} = \frac{\hbar^2}{2m_n} \left(\frac{2\pi}{\lambda} \right)^2 = \frac{81.81 \text{ meV} \cdot \text{\AA}^2}{\lambda^2}$$

$$E = k_B T = (0.08617 \text{ meV} \cdot \text{K}^{-1}) T$$



$\lambda \sim$ interatomic spacing $\rightarrow E \sim$ excitations in condensed matter

	Energy (meV)	Temperature (K)	Wavelength (\AA)
Cold	0.1 – 10	1 – 120	4 – 30
Thermal	5 – 100	60 – 1000	1 – 4
Hot	100 – 500	1000 – 6000	0.4 – 1

The Basic Experiment:



Incident Beam:

- monochromatic
- “white”
- “pink”

Scattered Beam:

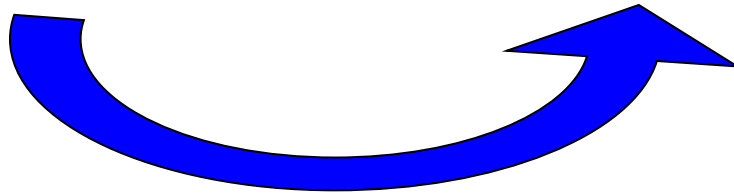
- Resolve its energy
- Don't resolve its energy
- Filter its energy

Fermi's Golden Rule within the 1st Born Approximation

$$W = 2\pi / \hbar \quad |\langle f | V | i \rangle|^2 \rho(E_f)$$



$$\delta\sigma = W / \Phi = (m/2\pi\hbar^2)^2 k_f / k_i |\langle f | V | i \rangle|^2 \delta\Omega$$



$$\delta^2\sigma / \delta\Omega \delta E_f = k_f / k_i \sigma_{\text{coh}} / 4\pi N S_{\text{coh}}(\mathbf{Q}, \omega)$$

$$+ k_f / k_i \sigma_{\text{incoh}} / 4\pi N S_{\text{incoh}}(\mathbf{Q}, \omega)$$

Nuclear correlation functions

Pair correlation function

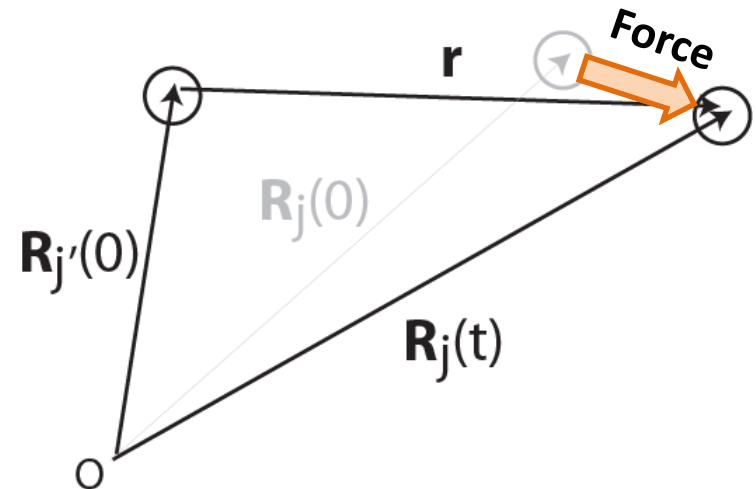
$$G(\mathbf{r}, t) = \frac{1}{N} \int \sum_{jj'} \delta(\mathbf{r}' - \mathbf{R}_{j'}(0)) \delta(\mathbf{r}' + \mathbf{r} - \mathbf{R}_j(t)) d\mathbf{r}'$$

Intermediate function

$$I(\mathbf{Q}, t) = \int G(\mathbf{r}, t) e^{i\mathbf{Q} \cdot \mathbf{r}} d\mathbf{r} = \frac{1}{N} \sum_{jj'} \exp(-i\mathbf{Q} \cdot \mathbf{R}_{j'}(0)) \exp(i\mathbf{Q} \cdot \mathbf{R}_j(t))$$

Scattering function

$$S(\mathbf{Q}, \omega) = \frac{1}{2\pi\hbar} \int I(\mathbf{Q}, t) e^{-i\omega t} dt$$



Differential scattering cross-section

$$\frac{d^2 \sigma}{d\Omega dE_f} = \frac{\sigma_{scat}}{4\pi} \frac{k_f}{k_i} NS(\mathbf{Q}, \omega)$$

Nuclear (lattice) excitations

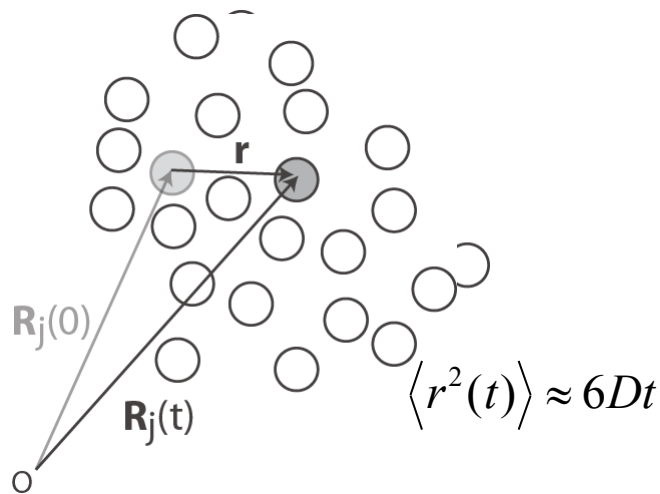
Neutron scattering measures simultaneously the wavevector and energy of **collective excitations** \rightarrow dispersion relation, $\omega(\mathbf{q})$

In addition, **local excitations** can of course be observed

- **Commonly studied excitations**
 - Phonons
 - Librations and vibrations in molecules
 - Diffusion
 - Collective modes in glasses and liquids
- **Excitations can tell us about**
 - Interatomic potentials & bonding
 - Phase transitions & critical phenomena (soft modes)
 - Fluid dynamics
 - Momentum distributions & superfluids (eg. He)
 - Interactions (eg. electron-phonon coupling)

Atomic diffusion

For long times compared to the collision time, atom diffuses

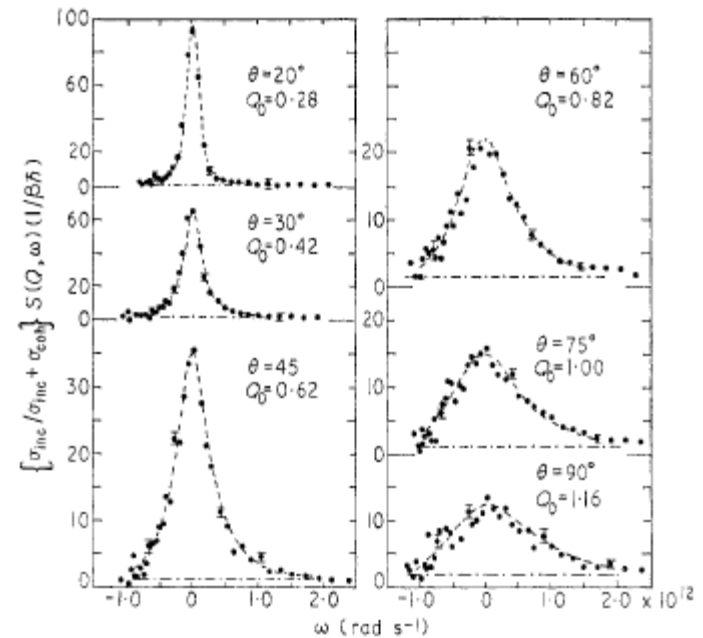


Auto-correlation function

$$G_s(r, t) = \left\{ 6\pi \langle r^2(t) \rangle \right\}^{-3/2} \exp\left(-\frac{r^2}{6 \langle r^2(t) \rangle} \right)$$

$$S(Q, \omega) = \frac{1}{\pi \hbar} \exp\left(\frac{\hbar \omega}{2k_B T} \right) \frac{DQ^2}{\omega^2 + (DQ^2)^2}$$

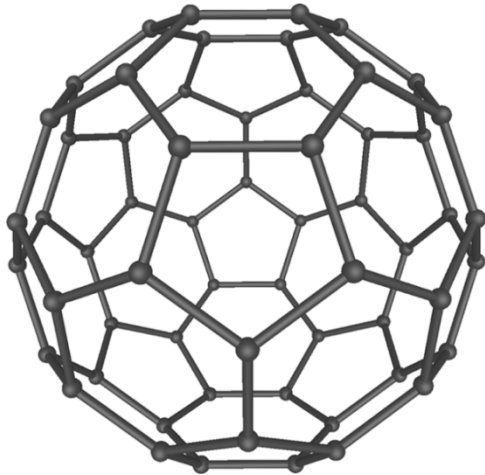
Liquid Na



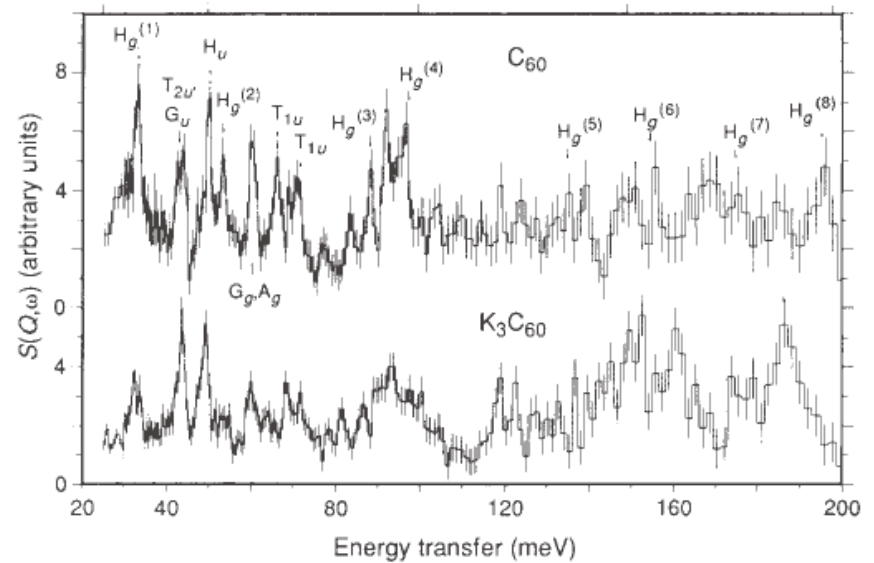
Cocking, *J. Phys. C* 2, 2047 (1969)..

Molecular vibrations

- Large molecule, many normal modes
- Harmonic vibrations can determine interatomic potentials



C60 molecule

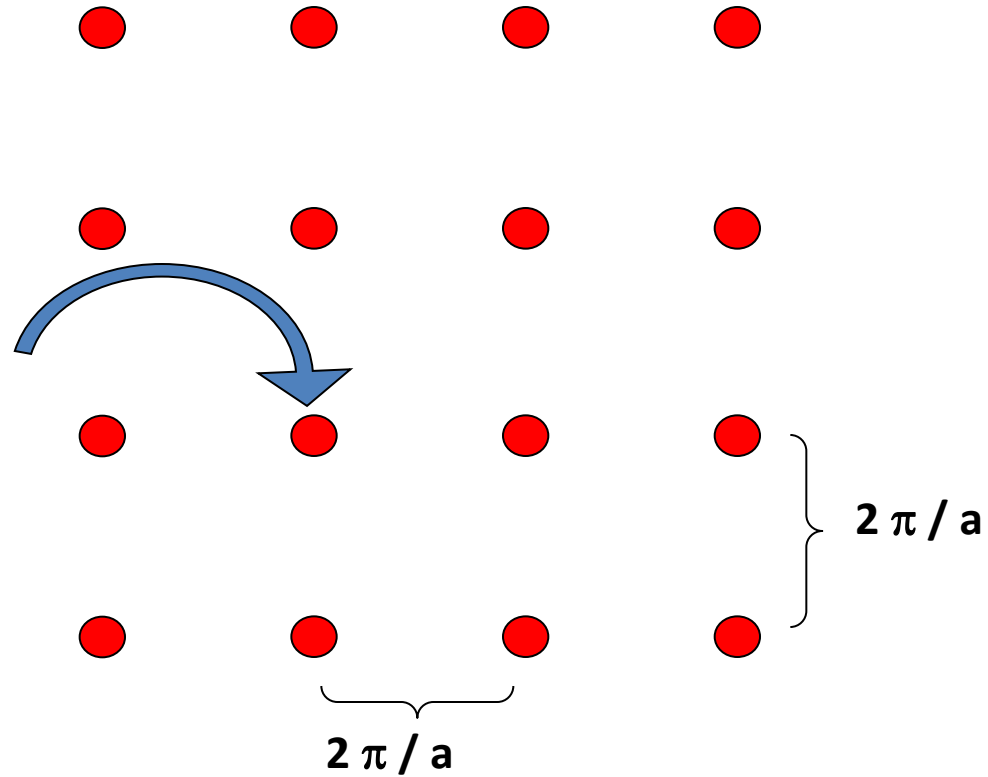


Prassides *et al.*, *Nature* **354**, 462 (1991).

Mapping Momentum – Energy (Q-E) space

Origin of reciprocal space;

Remains fixed for any sample rotation

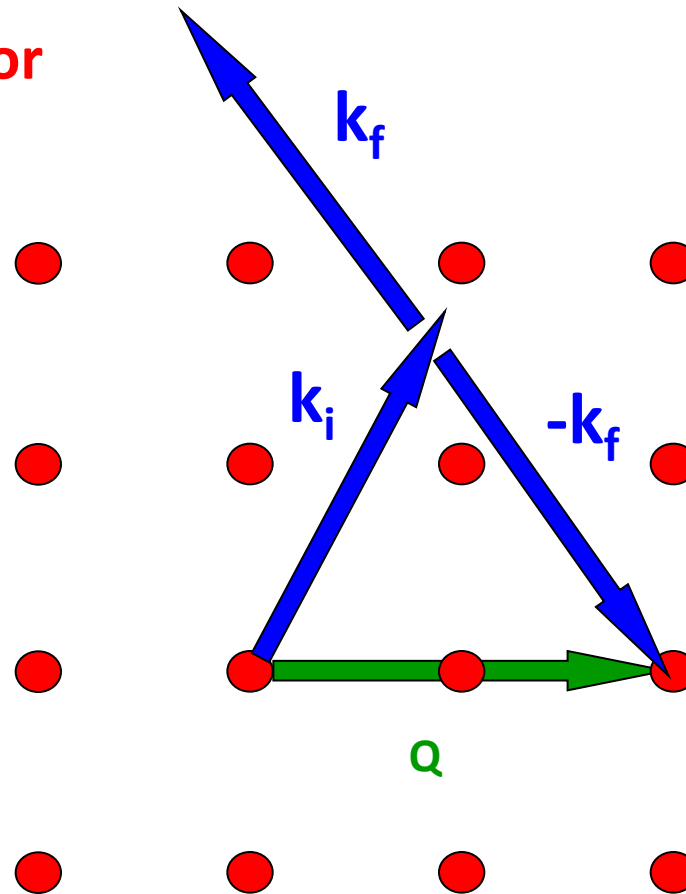


Bragg diffraction:

Constructive Interference

Q = Reciprocal Lattice Vector

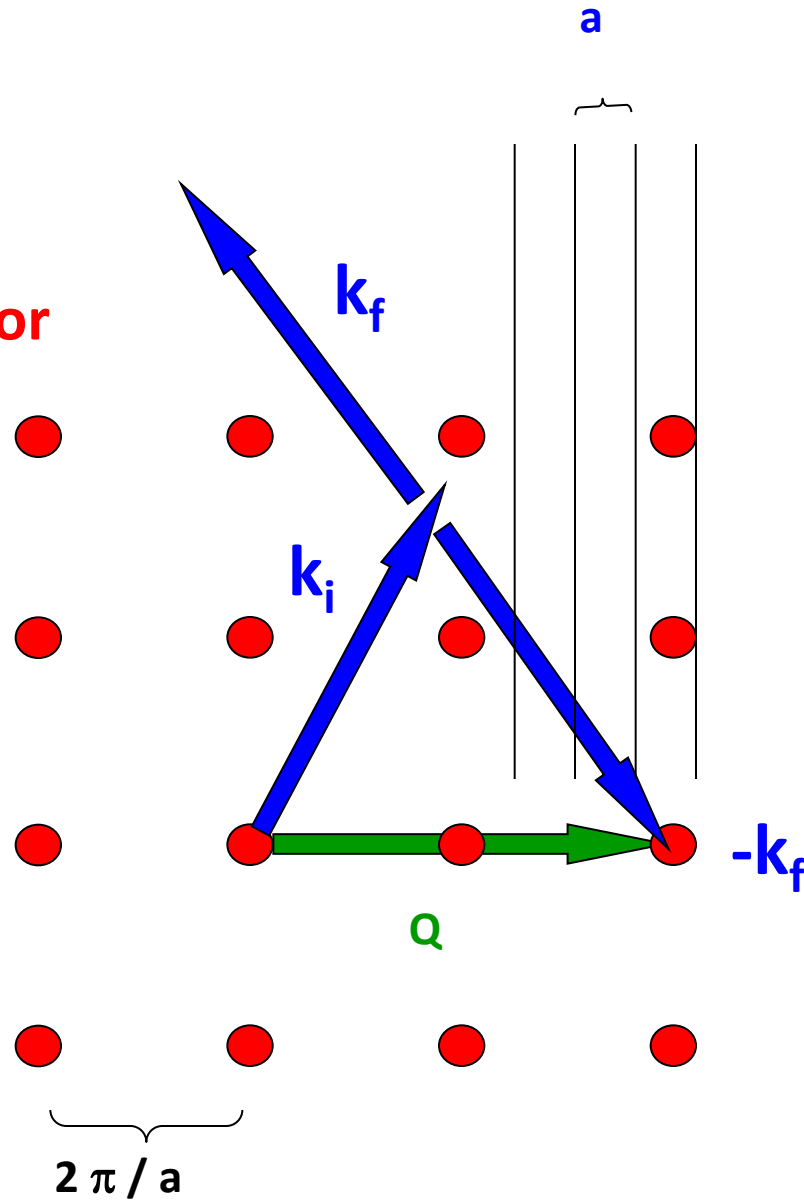
Elastic scattering : $| \mathbf{k}_i | = | \mathbf{k}_f |$



Bragg diffraction:

Constructive Interference

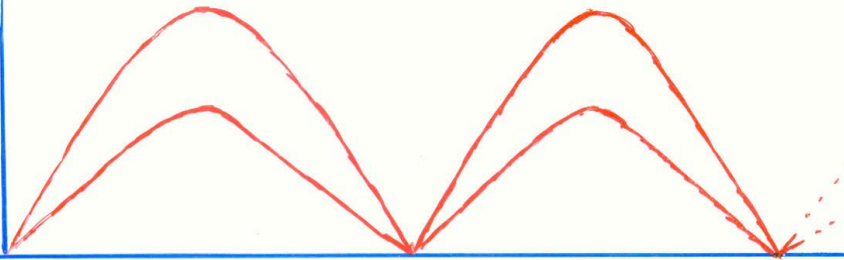
Q = Reciprocal Lattice Vector



Elastic scattering : $|k_i| = |k_f|$

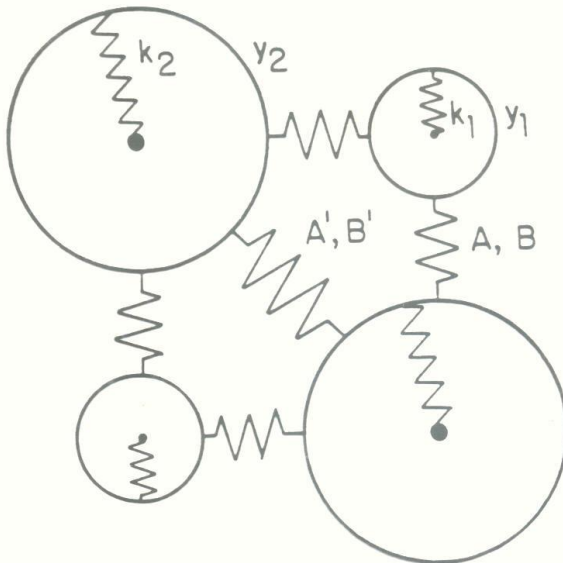
Elementary Excitations in Solids

ENERGY



MOMENTUM, \bar{Q}

- Lattice Vibrations (Phonons)
- Spin Fluctuations (Magnons)



Energy vs Momentum

- Forces which bind atoms together in solids

Phonons

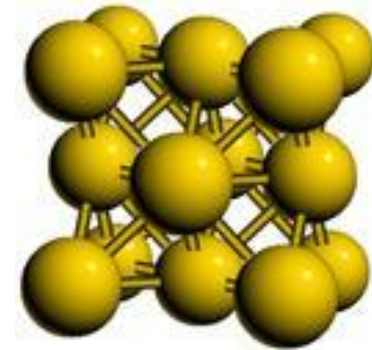
- Normal modes in periodic crystal → wavevector

$$\mathbf{u}(l,t) = \frac{1}{\sqrt{NM}} \sum_{j\mathbf{q}} \boldsymbol{\varepsilon}_j(\mathbf{q}) \exp(i\mathbf{q} \cdot \mathbf{l}) \hat{B}(\mathbf{q}j,t)$$

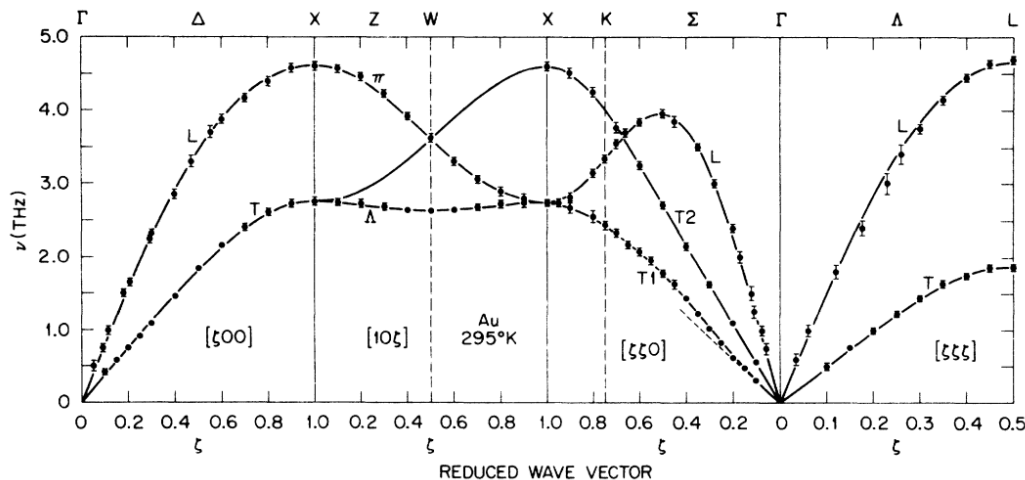
- Energy of phonon depends on \mathbf{q} and polarization



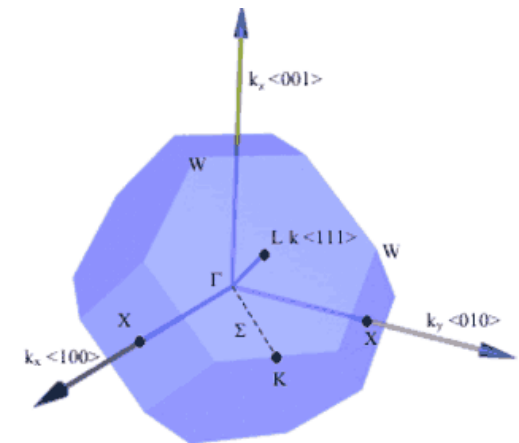
Longitudinal mode



FCC structure



Lynn, et al., *Phys. Rev. B* **8**, 3493 (1973).

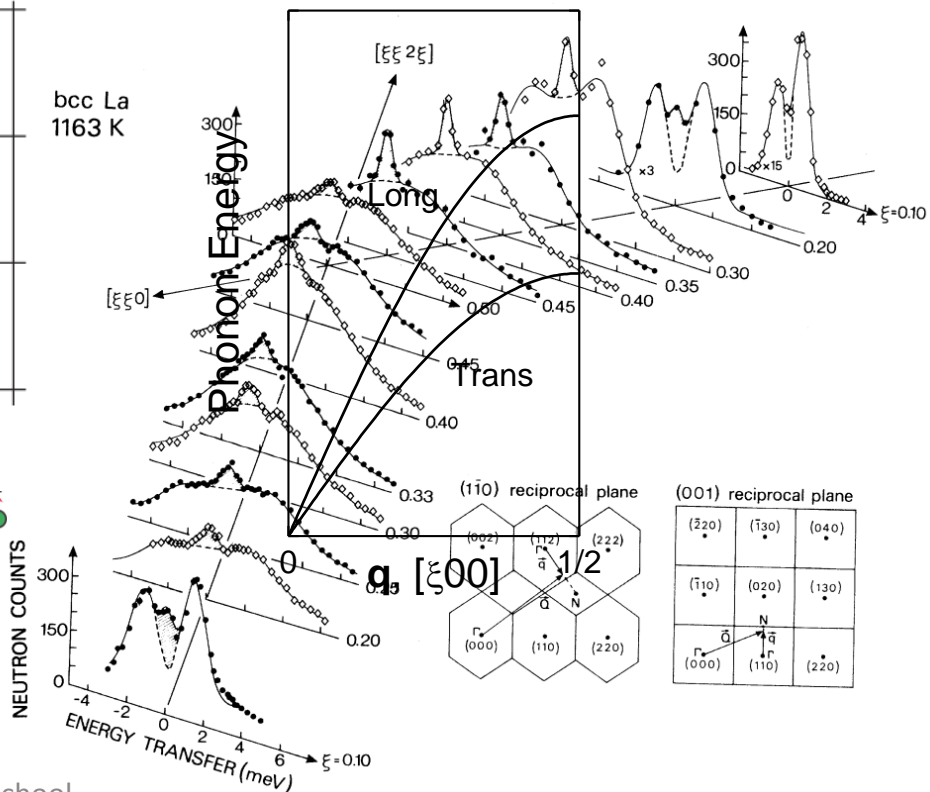
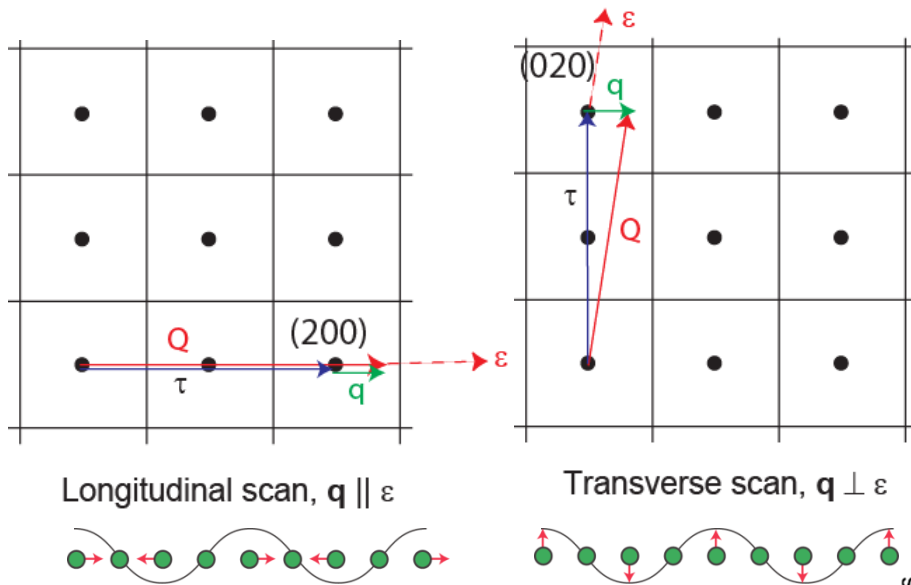


FCC Brillouin zone

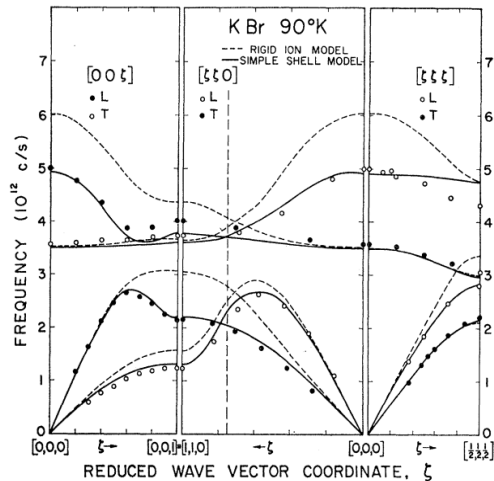
Phonon intensities

$$S_{1+}(\mathbf{Q}, \omega) = \frac{1}{2NM} e^{-Q^2 \langle u^2 \rangle} \sum_{j\mathbf{q}} \frac{|\mathbf{Q} \cdot \boldsymbol{\varepsilon}_j(\mathbf{q})|^2}{\omega_j(\mathbf{q})} (1 + n(\omega)) \delta(\mathbf{Q} - \mathbf{q} - \boldsymbol{\tau}) \delta(\omega - \omega_j(\mathbf{q}))$$

← Structure (polarization) factor



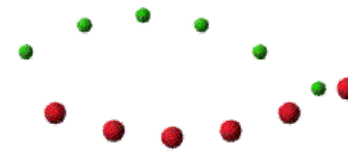
More complicated structures



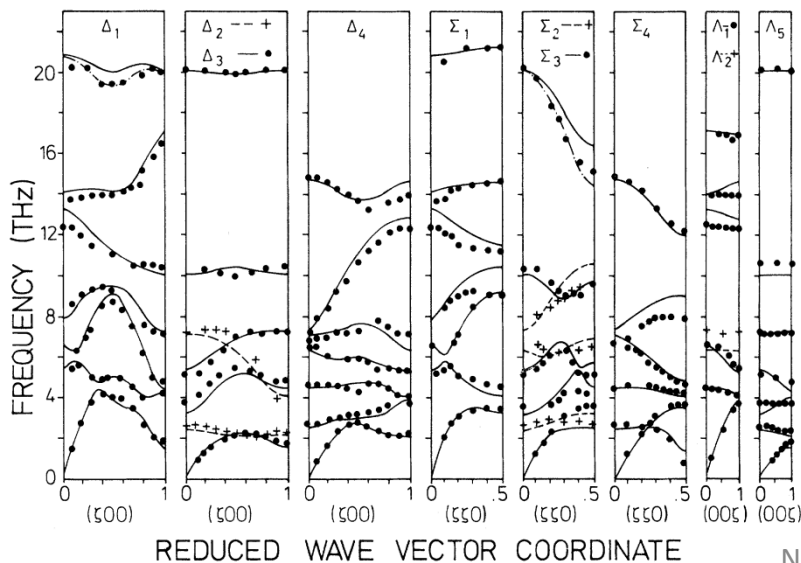
Woods, et al., *Phys. Rev.* **131**, 1025 (1963).



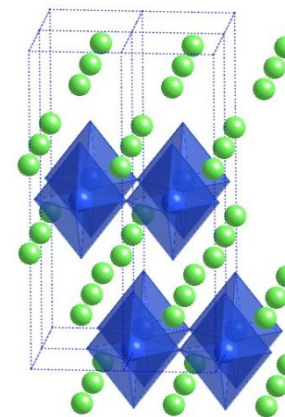
Acoustic phonon



Optical phonon



Chaplot, et al., *Phys. Rev. B* **52**, 7230(1995).



La_2CuO_4

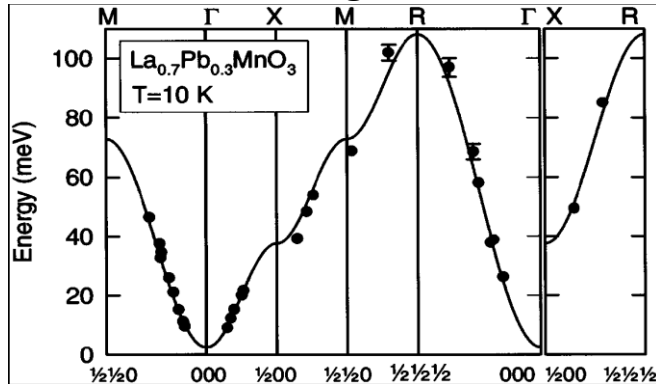
Spin excitations

- **Spin excitations**
 - Spin waves in ordered magnets
 - Paramagnetic & quantum spin fluctuations
 - Crystal-field & spin-orbit excitations

- **Magnetic inelastic scattering can tell us about**
 - Exchange interactions
 - Single-ion and exchange anisotropy (determine Hamiltonian)
 - Phase transitions & critical phenomena
 - Quantum critical scaling of magnetic fluctuations
 - Other electronic energy scales (eg. CF & SO)
 - Interactions (eg. spin-phonon coupling)

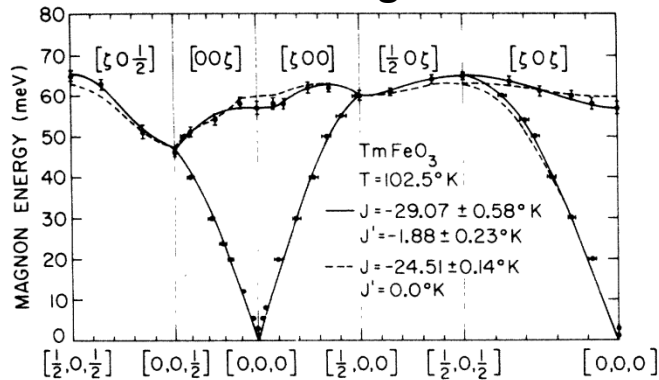
Spin waves

Ferromagnetic



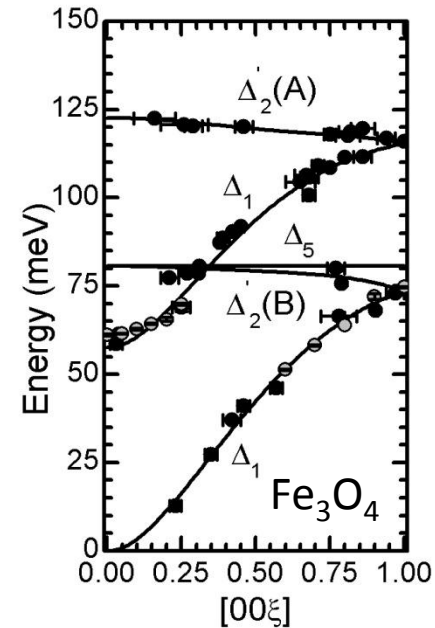
Perring *et al.*, *Phys. Rev. Lett.* **77**, 711 (1996).

Antiferromagnetic



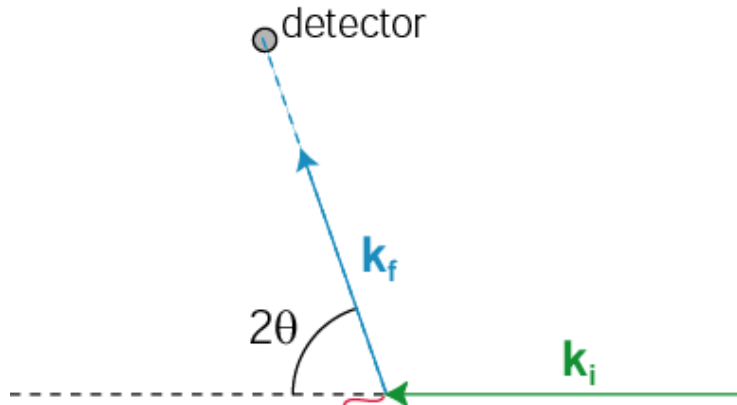
Shapiro *et al.*, *Phys. Rev. B* **10**, 2014 (1974).

Ferrimagnetic



McQueeney *et al.*, *Phys. Rev. Lett.* **99**, 246401 (2007).

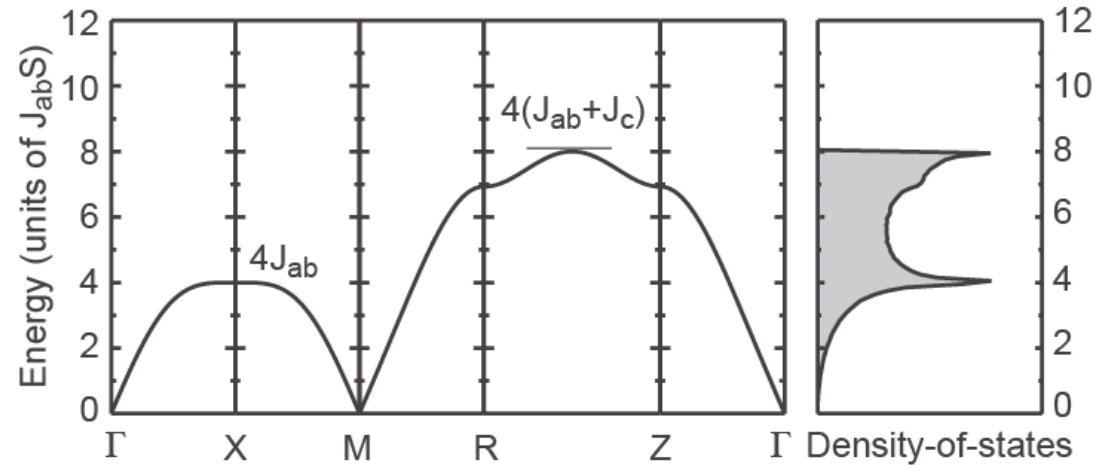
Scattering experiments



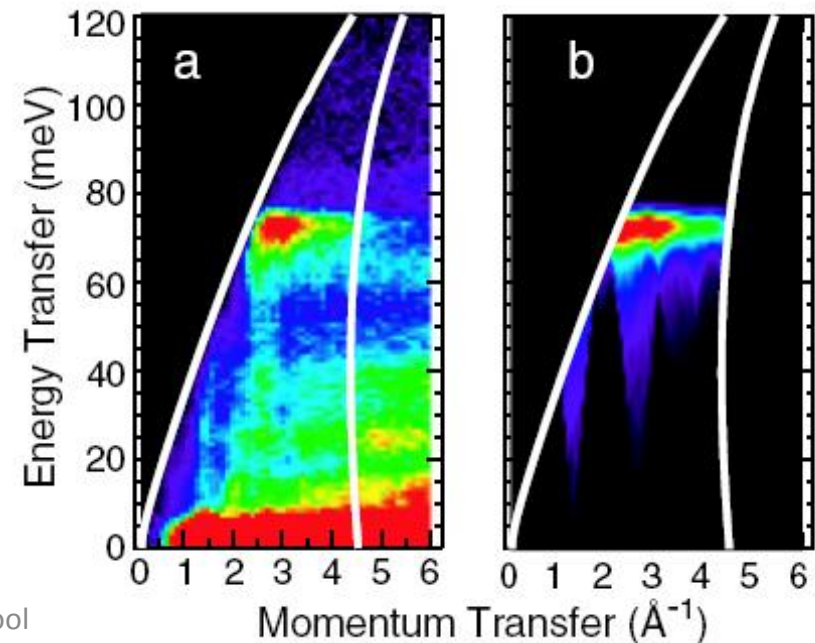
excitation $\hbar\omega = \frac{\hbar^2}{2m} (k_i^2 - k_f^2)$
 $\mathbf{Q} = \mathbf{k}_i - \mathbf{k}_f$

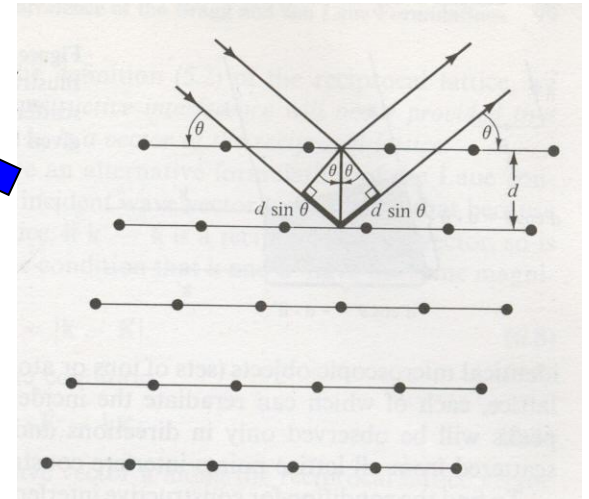
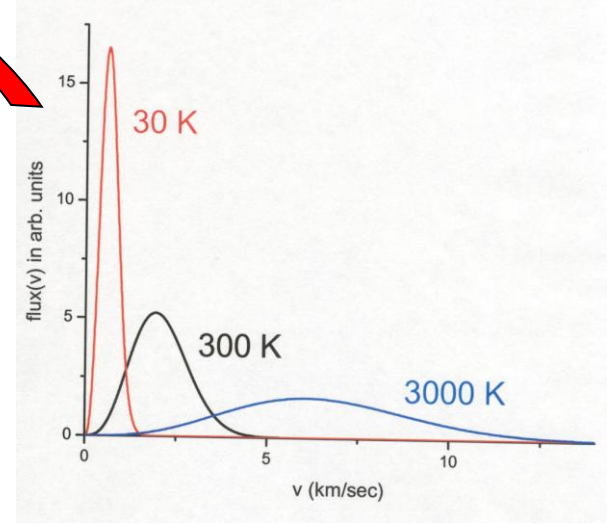
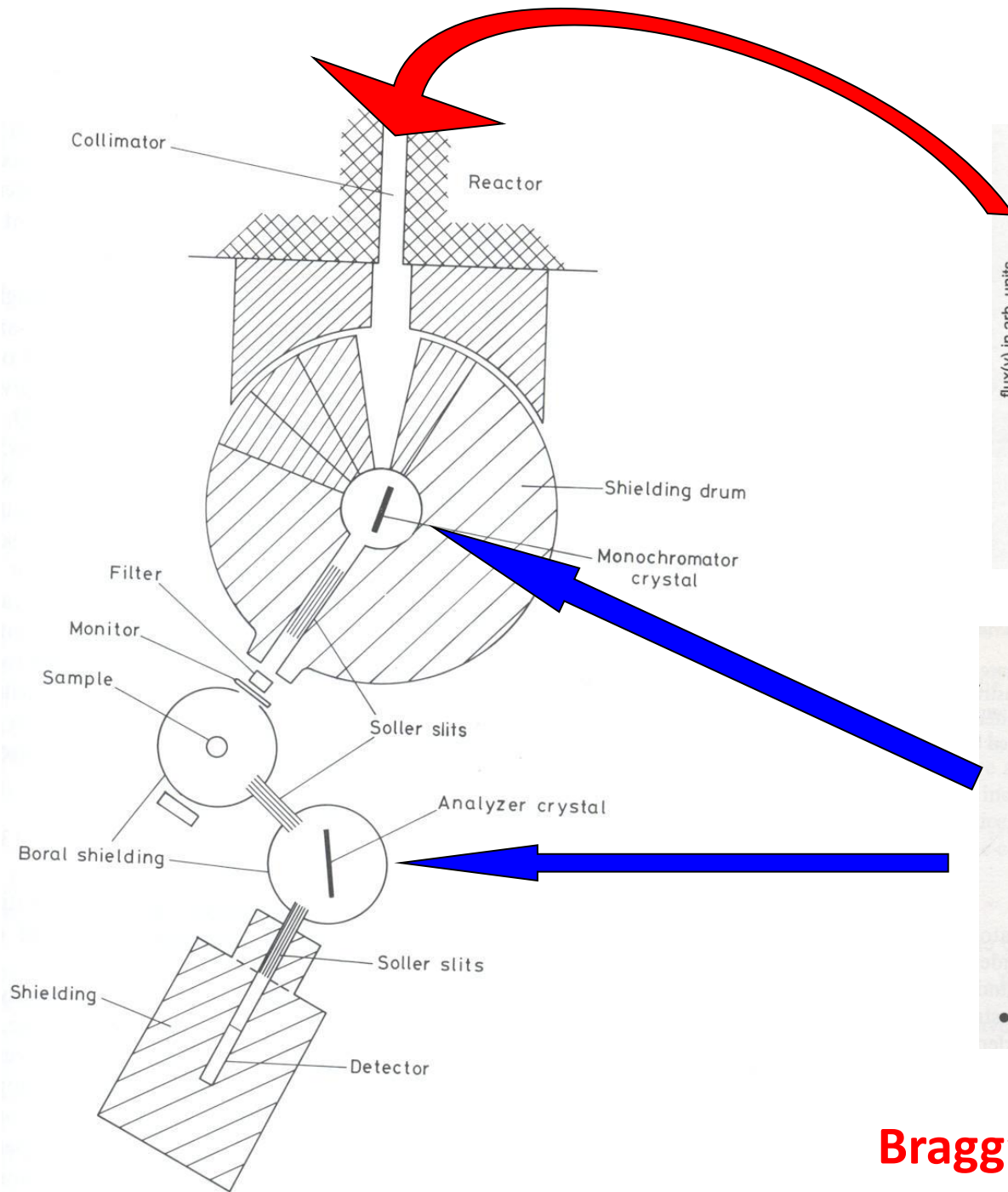
Instrument and sample (powder or single-crystal) determine how (\mathbf{Q}, ω) space is sampled

Single-crystal

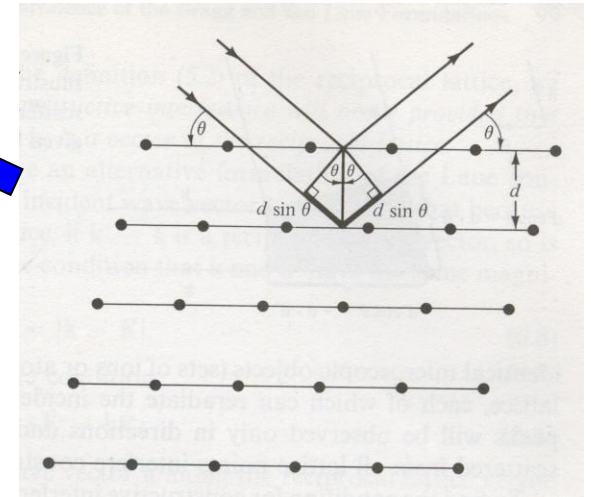
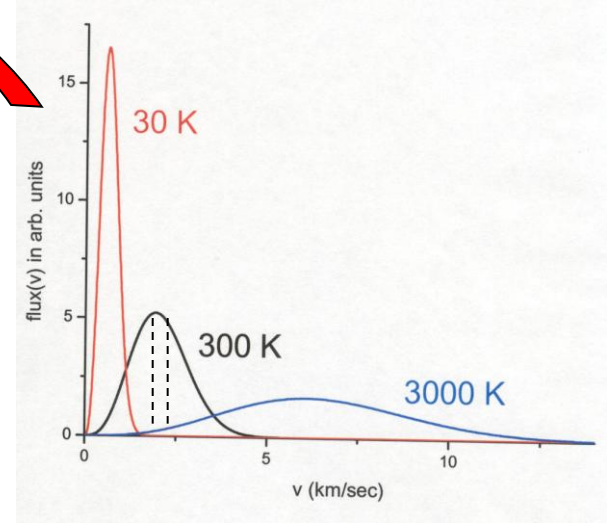
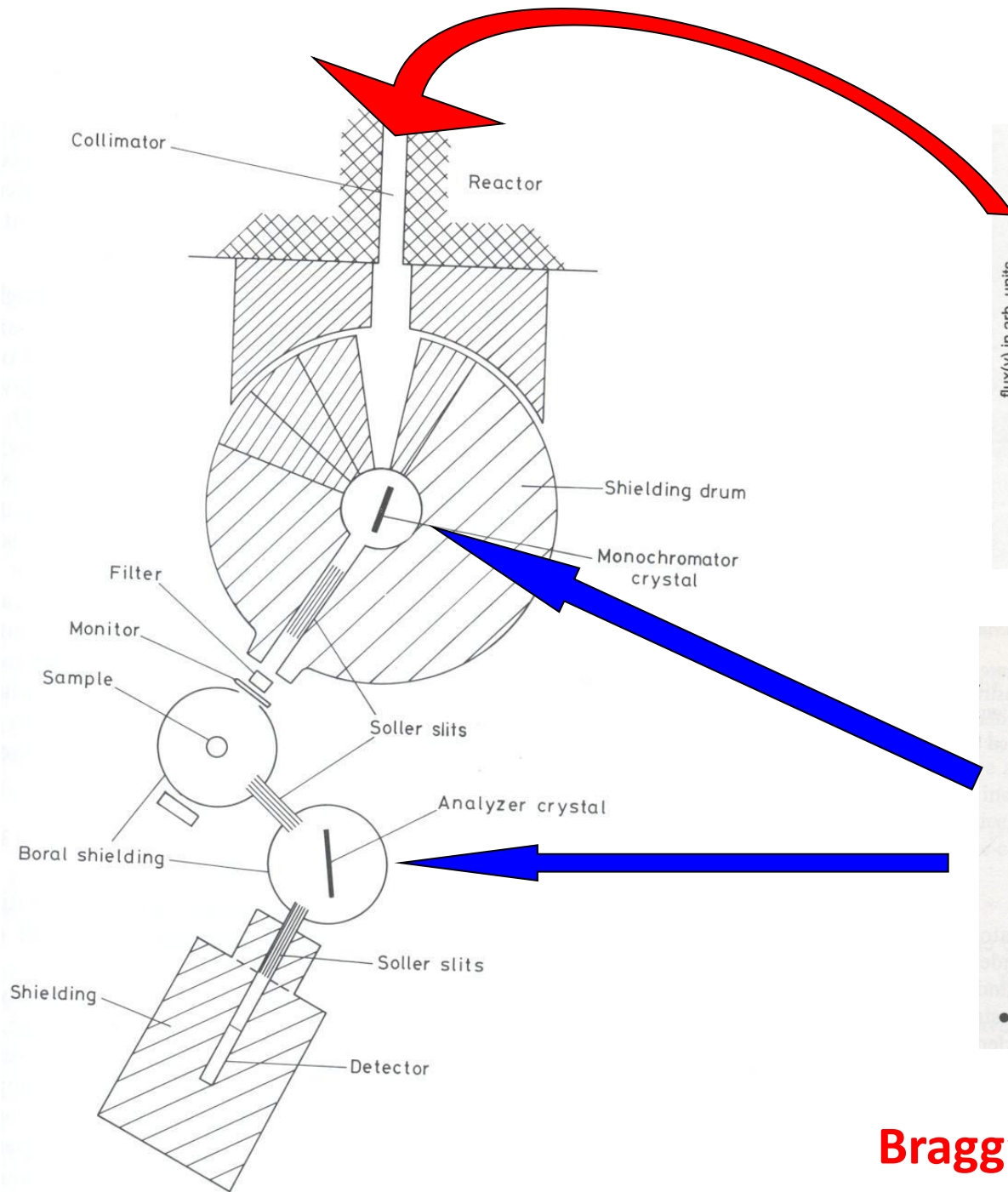


Powder $S(|\mathbf{Q}|, \omega)$





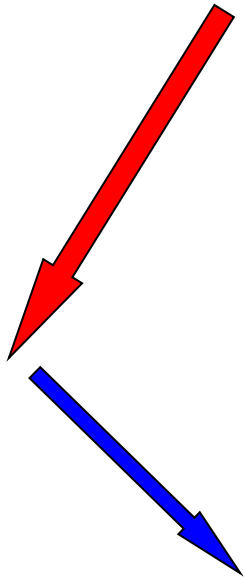
Bragg's Law: $n\lambda = 2d \sin(\theta)$



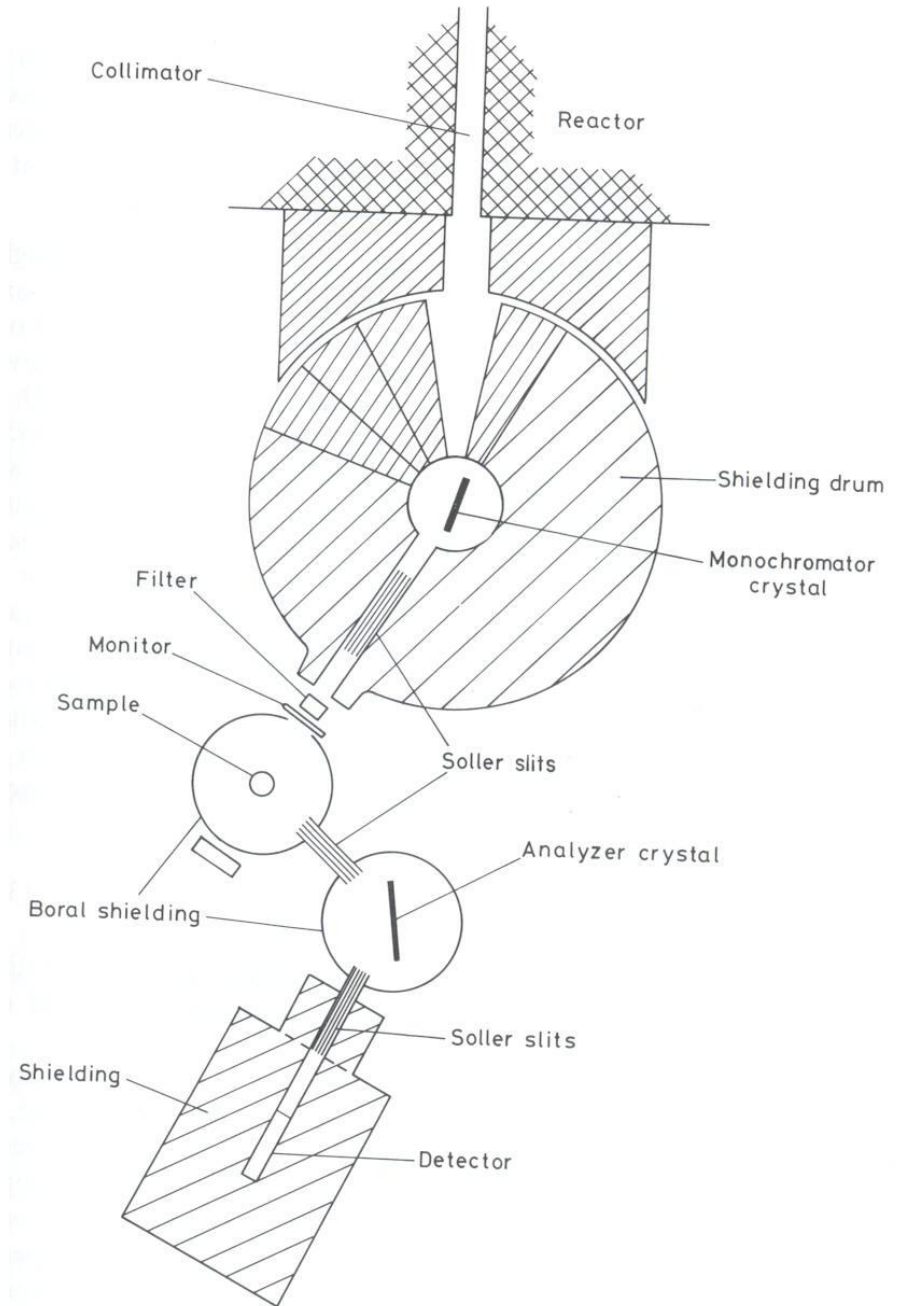
Bragg's Law: $n\lambda = 2d \sin(\theta)$

Brockhouse's Triple Axis Spectrometer

$$|\mathbf{k}_i| = 2\pi / \lambda_i$$

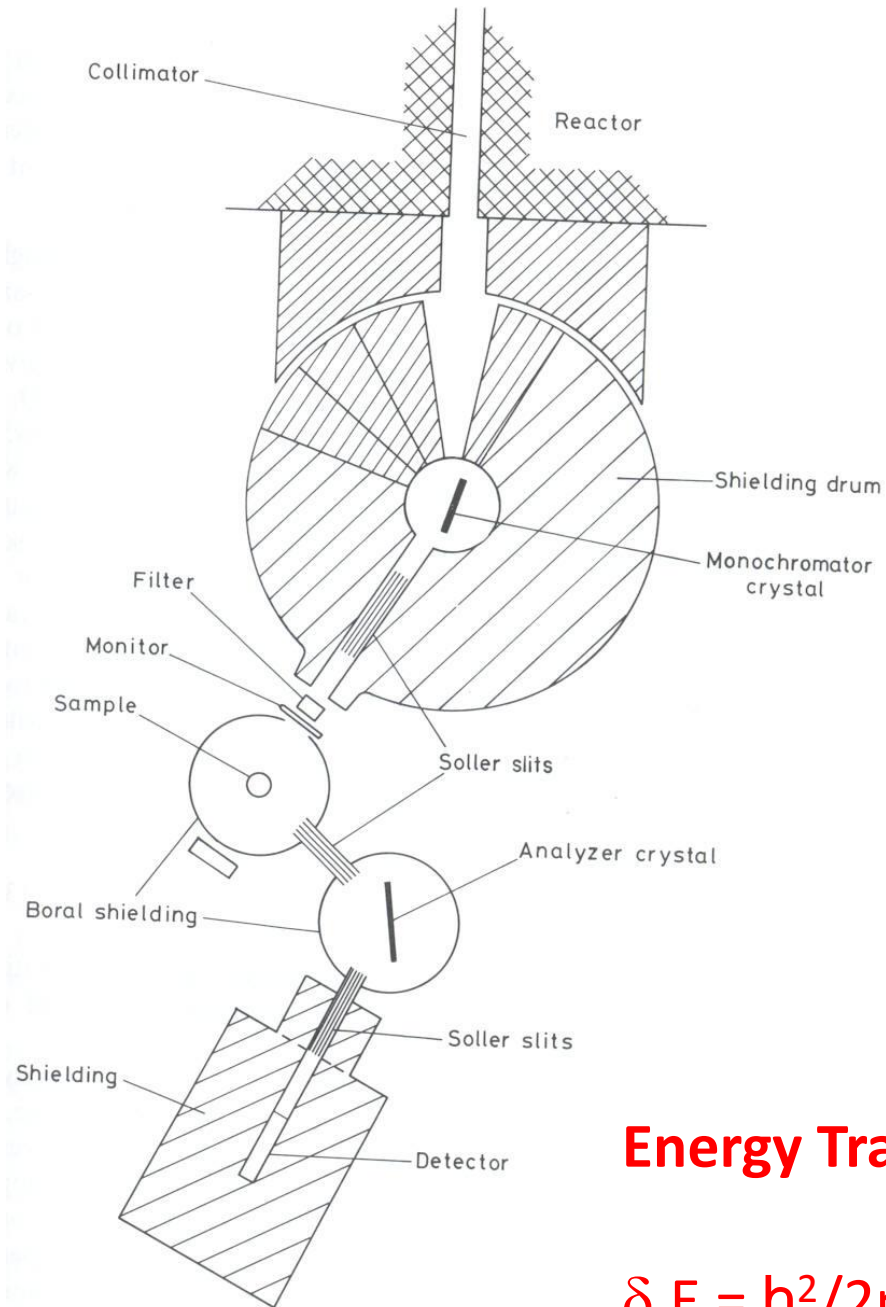
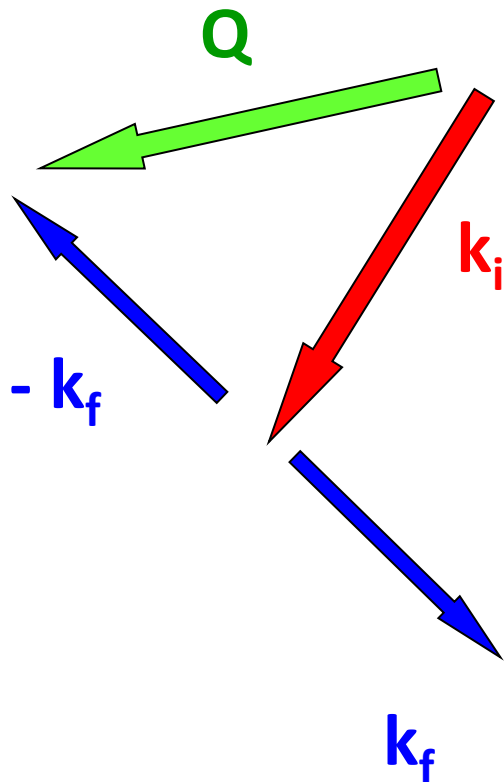


$$|\mathbf{k}_f| = 2\pi / \lambda_f$$



Momentum Transfer:

$$Q = k_i - k_f$$



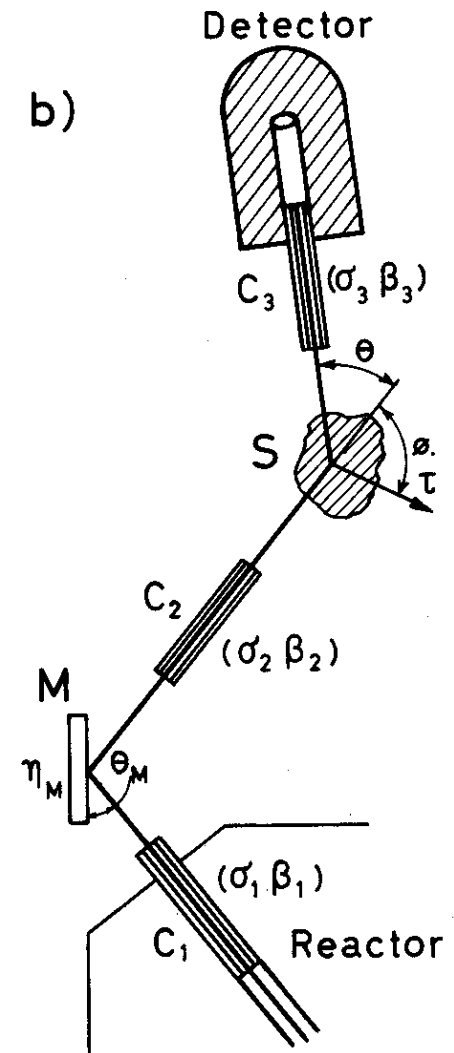
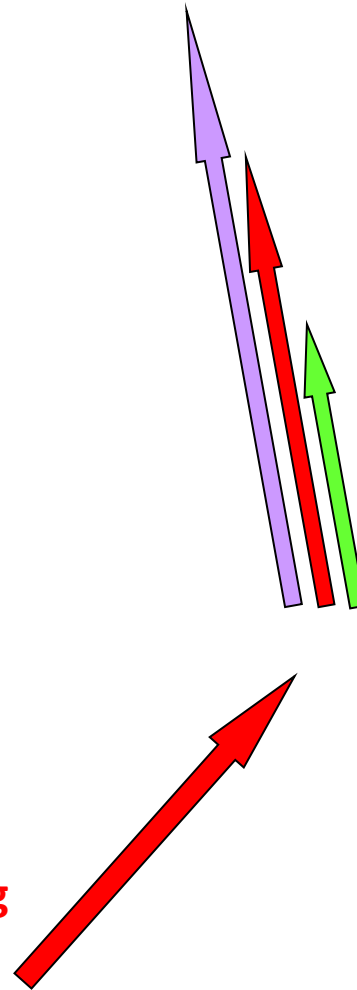
Energy Transfer:

$$\delta E = \frac{h^2}{2m} (k_i^2 - k_f^2)$$

Two Axis Spectrometer:

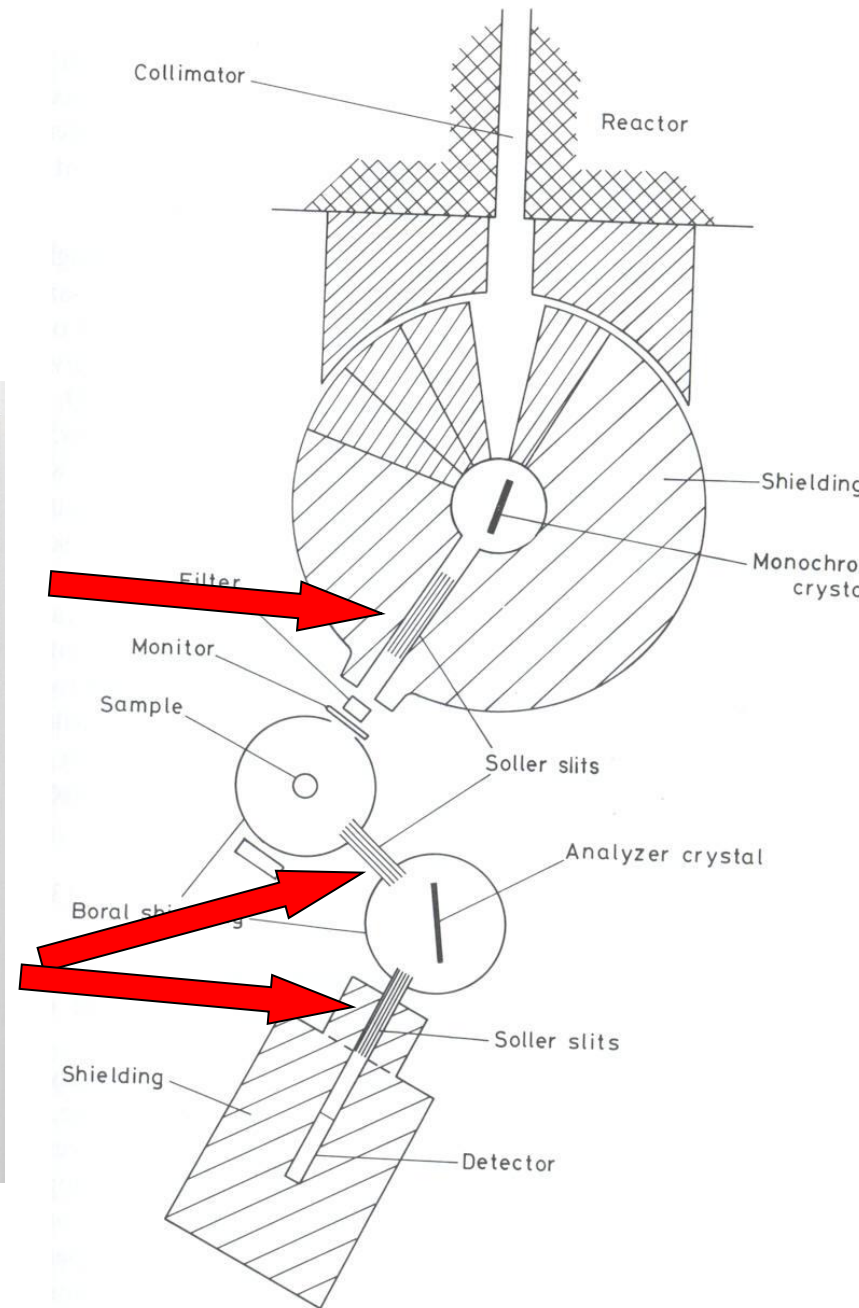
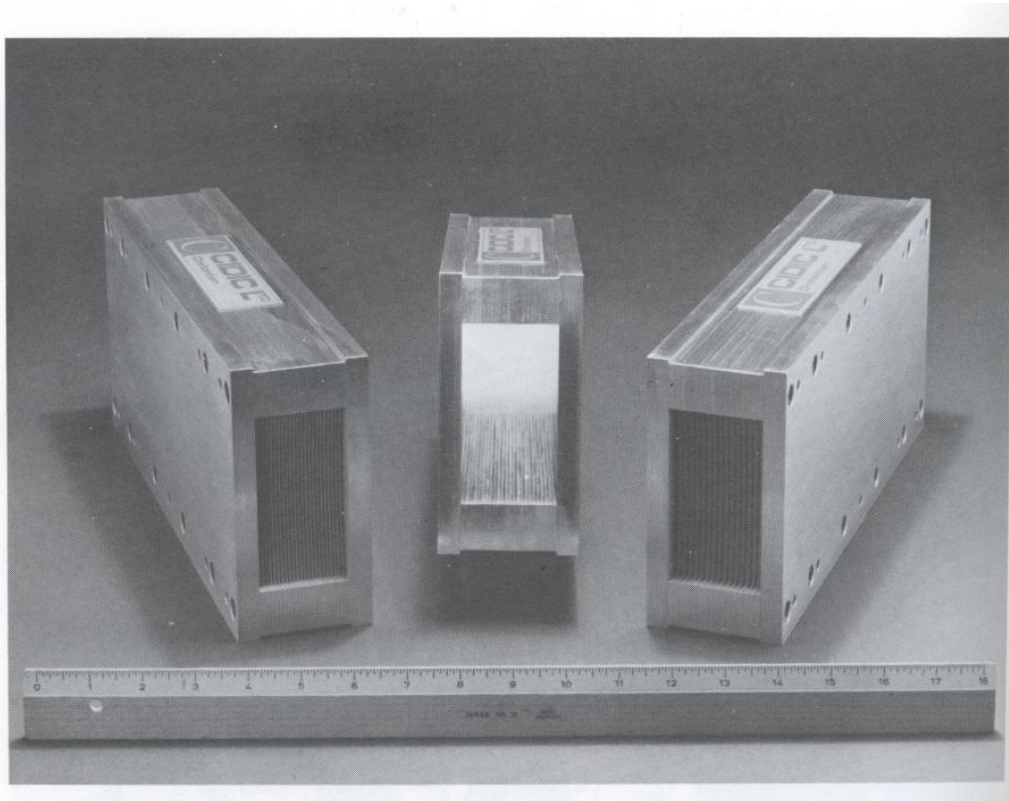
- 3-axis with analyser removed
- Powder diffractometer
- Small angle diffractometer
- Reflectometers

Diffractometers often employ working assumption that all scattering is *elastic*.



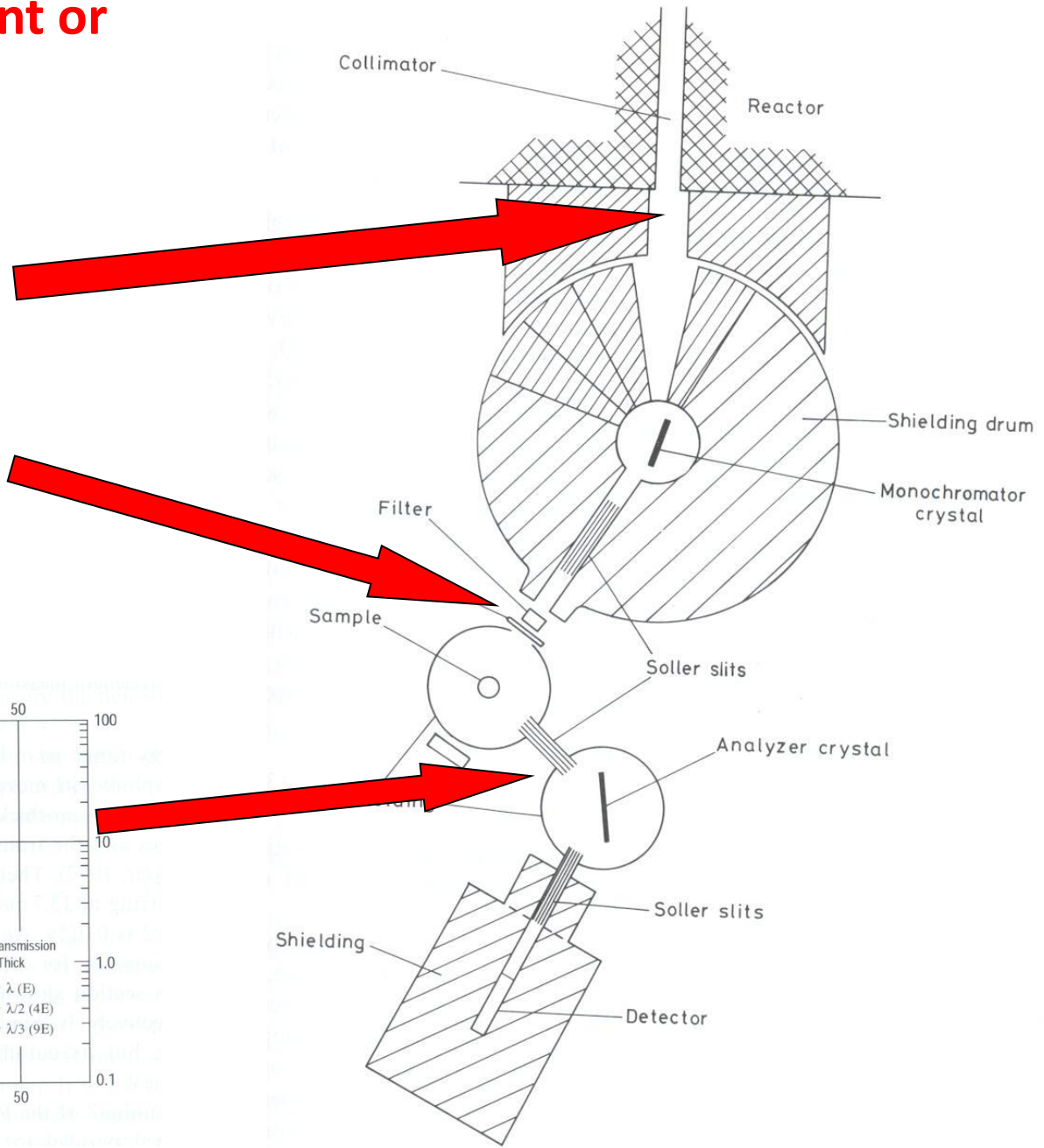
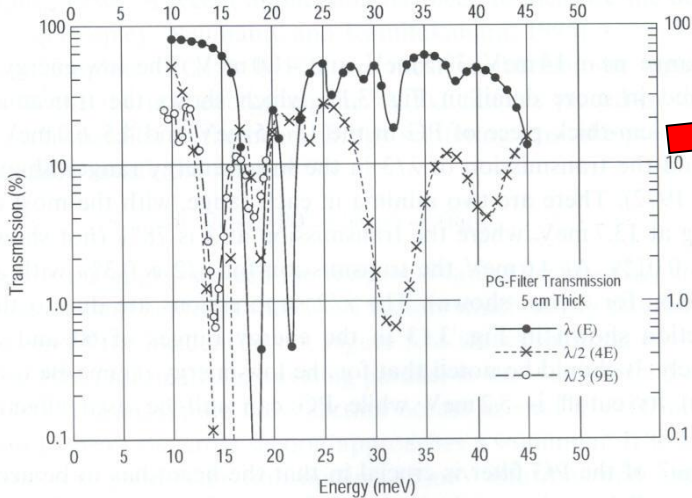
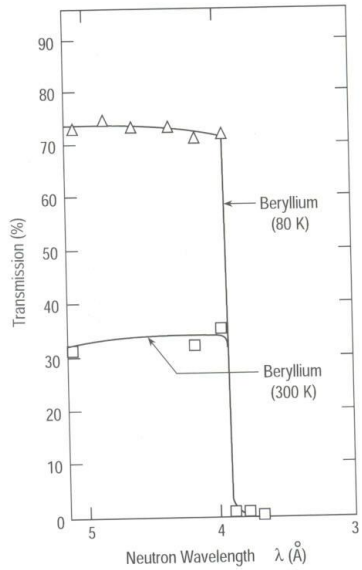
Soller Slits: Collimators

Define beam direction to
 $\pm 0.5, 0.75$ etc. degrees



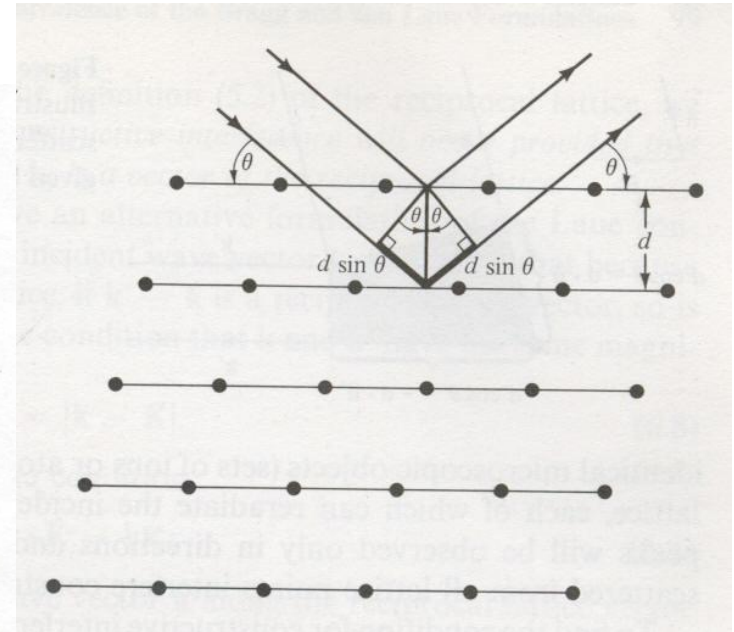
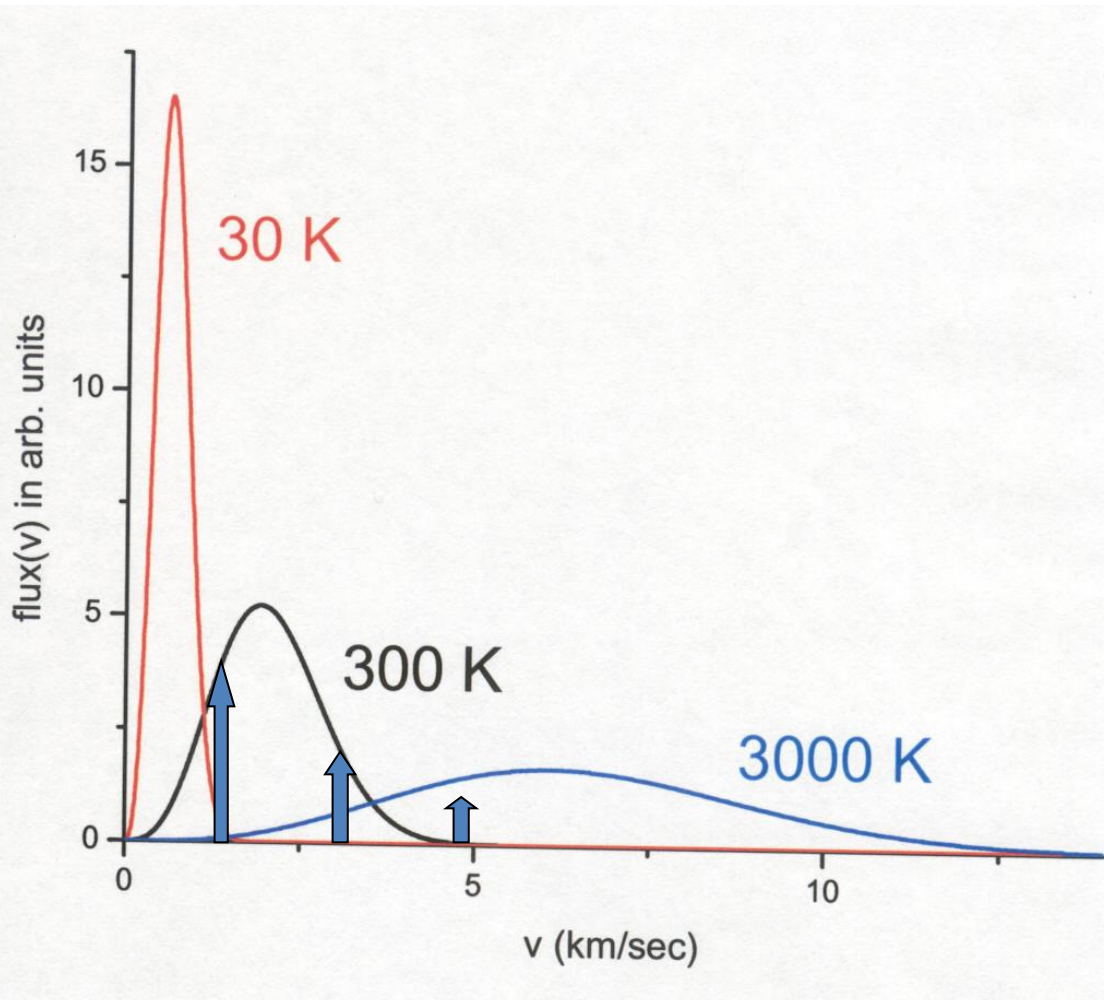
Filters:

Remove λ / n from incident or scattered beam, or both



Single crystal monochromators:

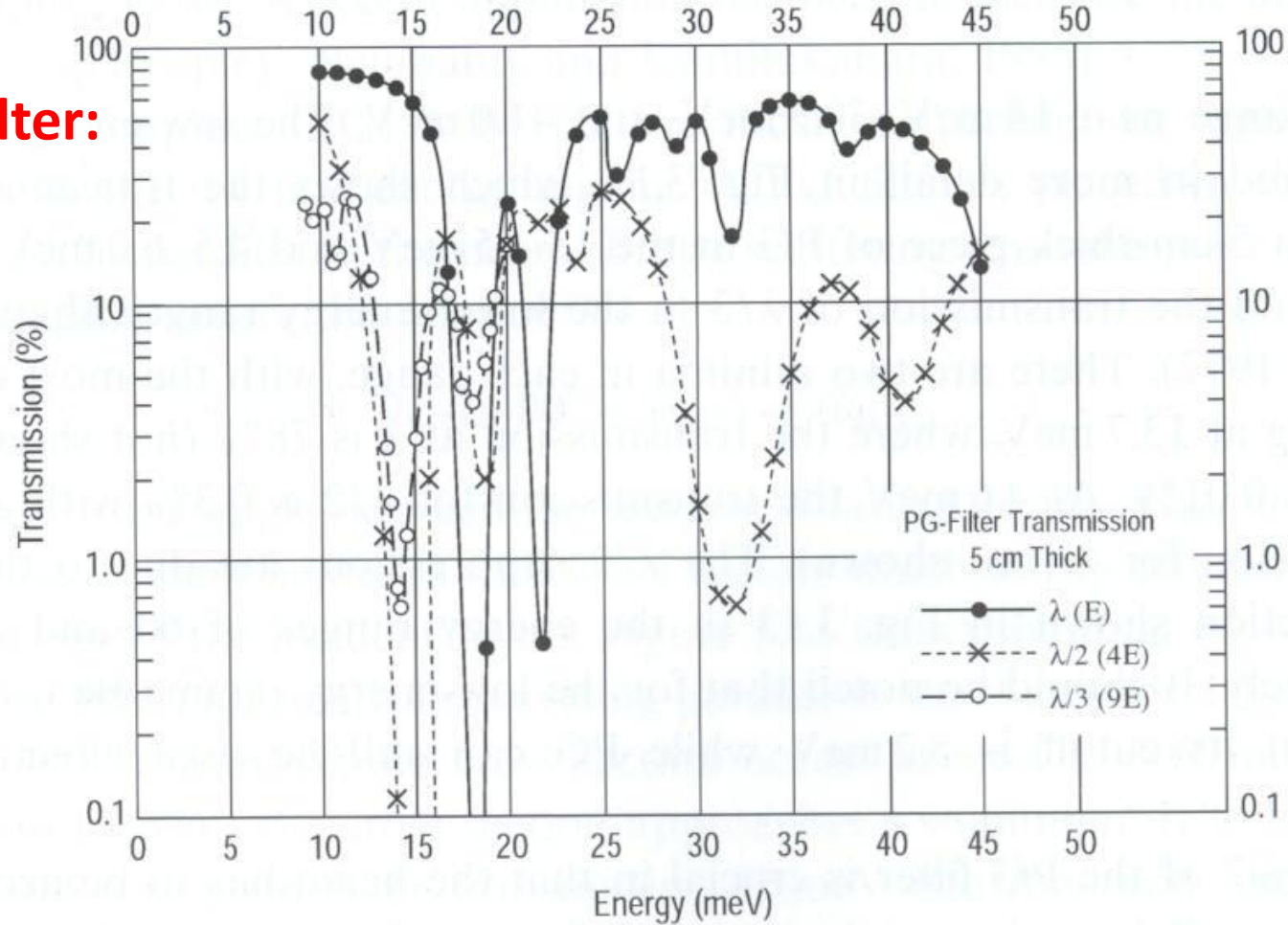
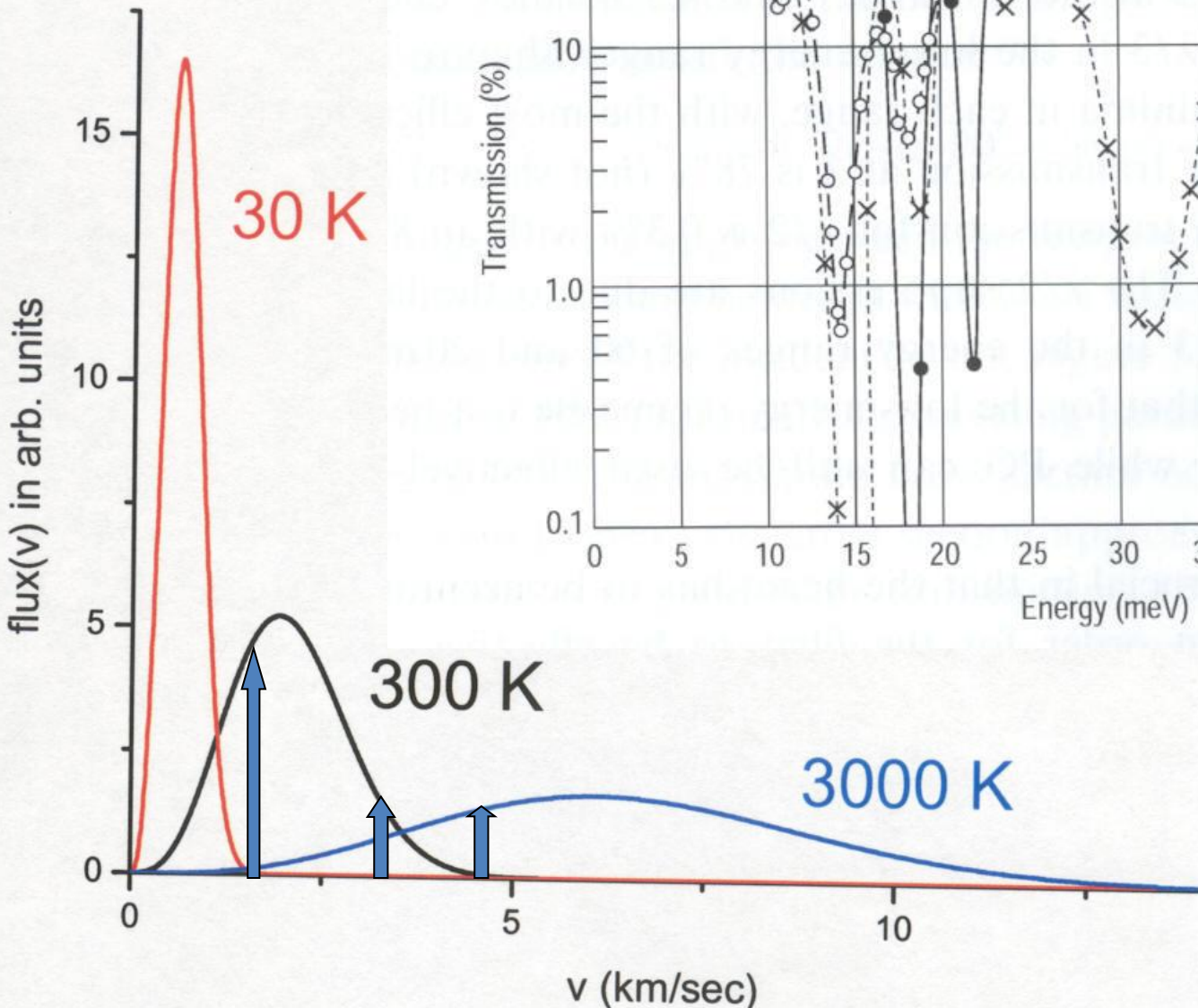
Bragg reflection and harmonic contamination



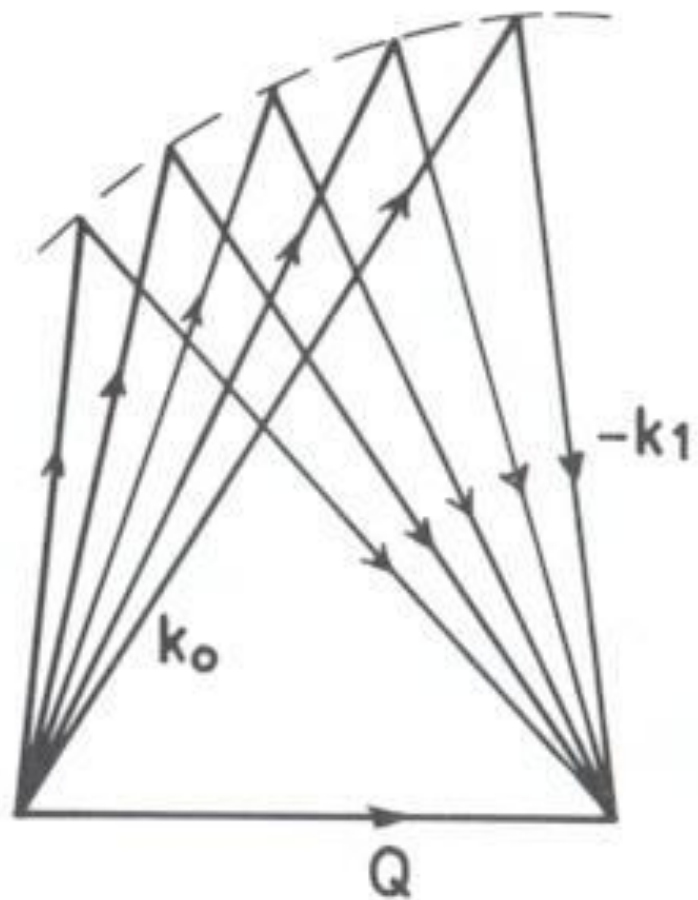
$$n\lambda = 2d \sin(\theta)$$

Get: $\lambda, \lambda/2, \lambda/3, \text{ etc.}$

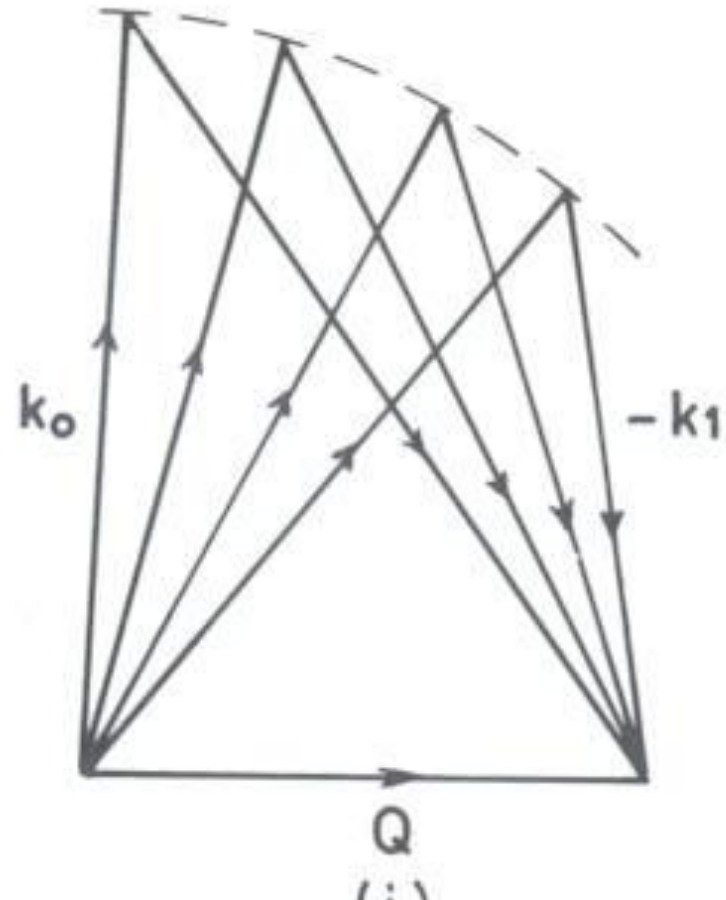
Pyrolytic graphite filter:



$E = 14.7 \text{ meV}$
 $\lambda = 2.37 \text{ \AA}$
 $v = 1.6 \text{ km/s}$
 $2 \times v = 3.2 \text{ km/s}$
 $3 \times v = 4.8 \text{ km/s}$



Constant k_f (ii)



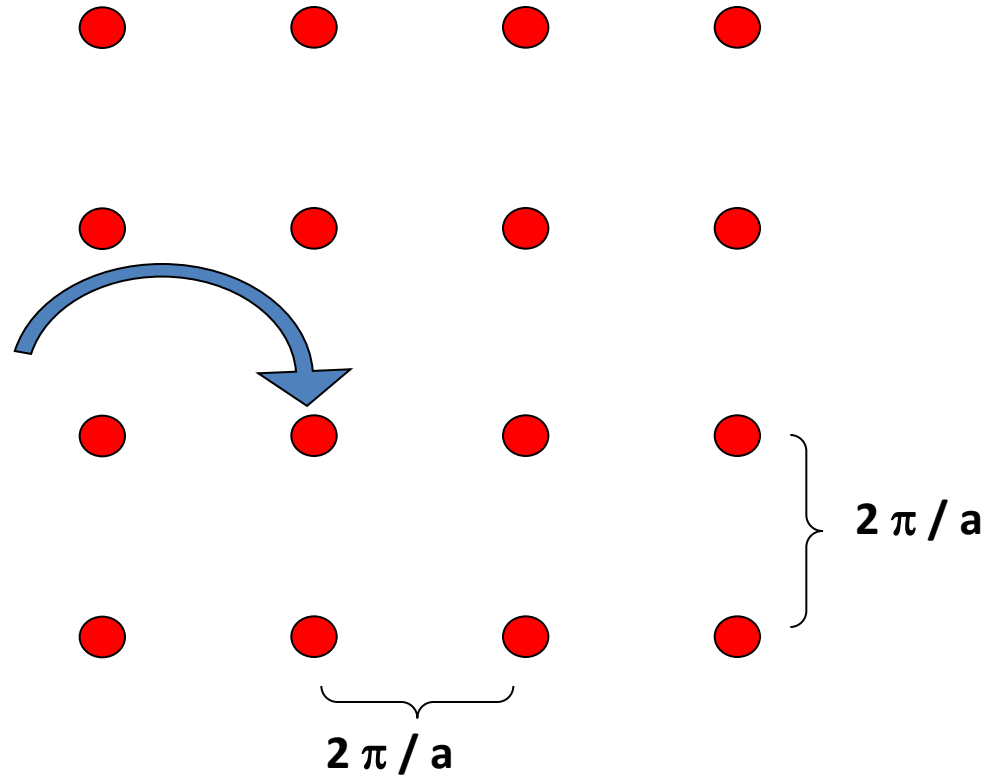
Constant k_i (i)

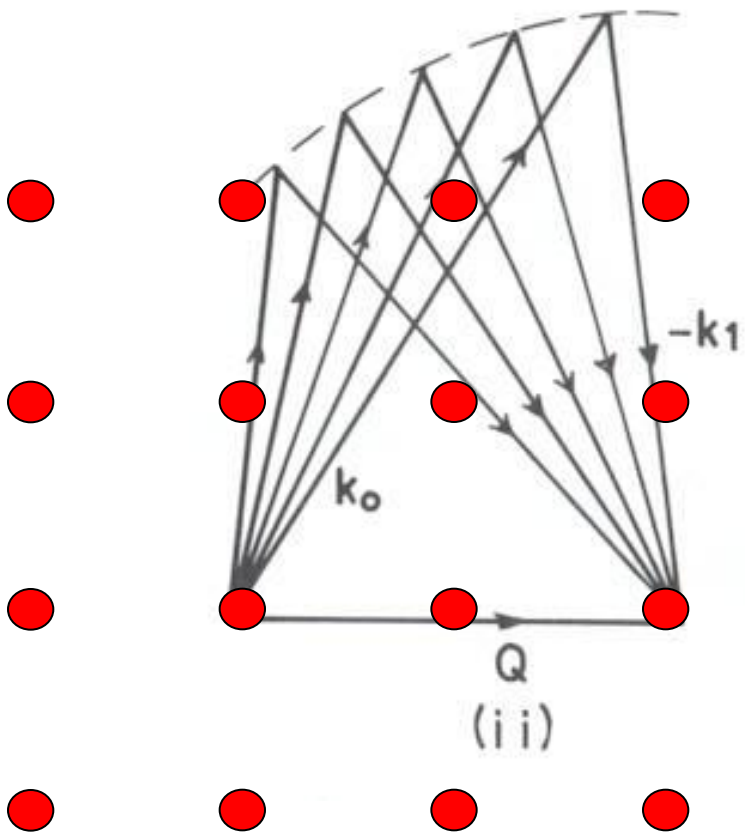
Two different ways of performing constant-Q scans

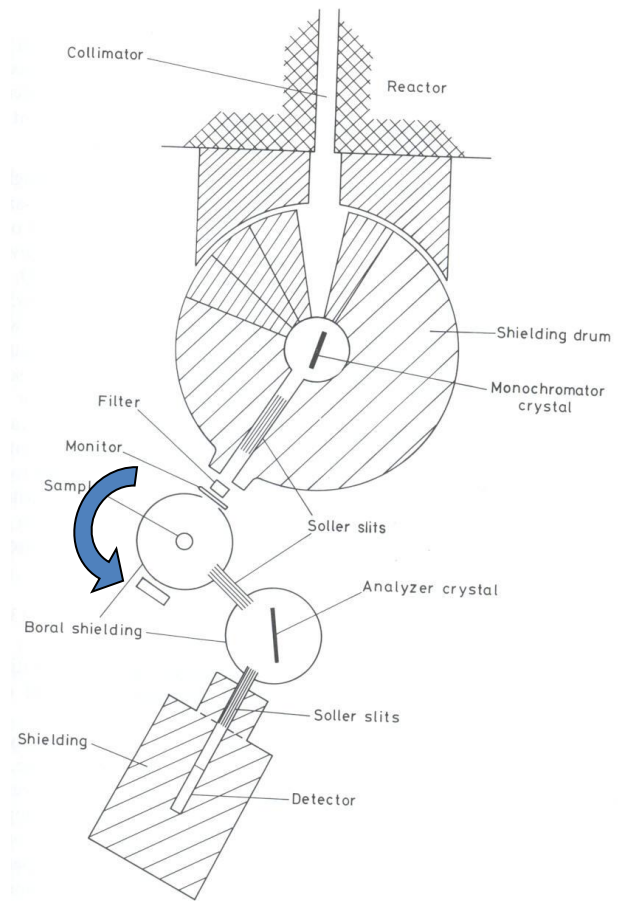
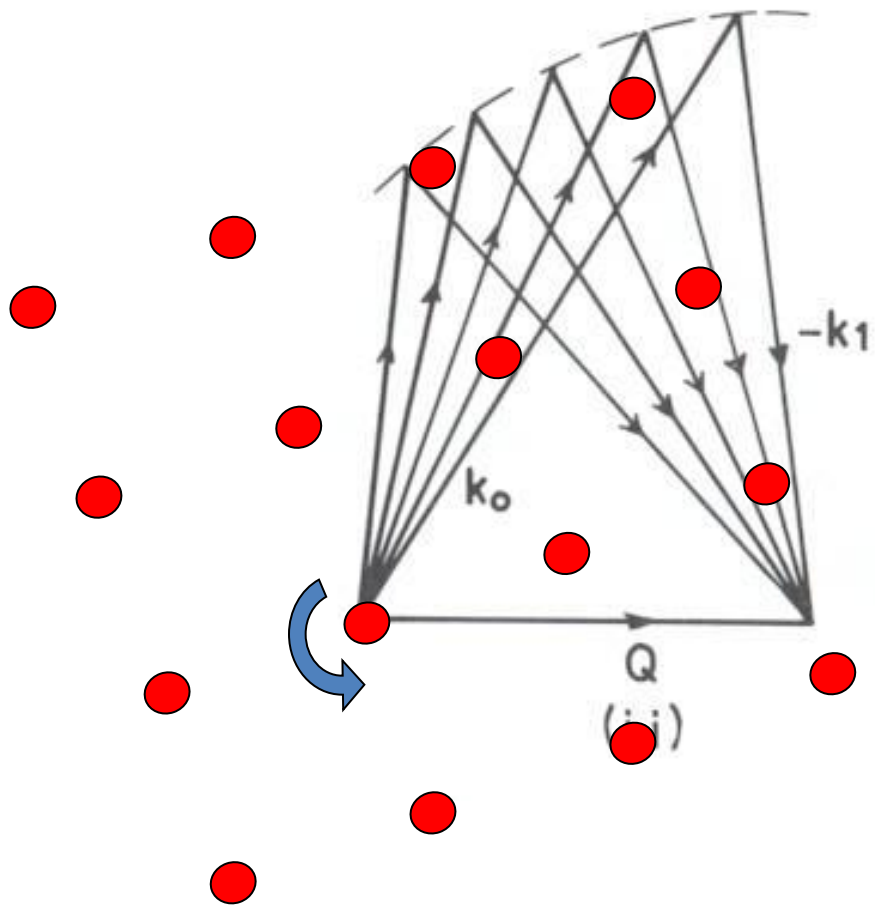
Mapping Momentum – Energy (Q-E) space

Origin of reciprocal space;

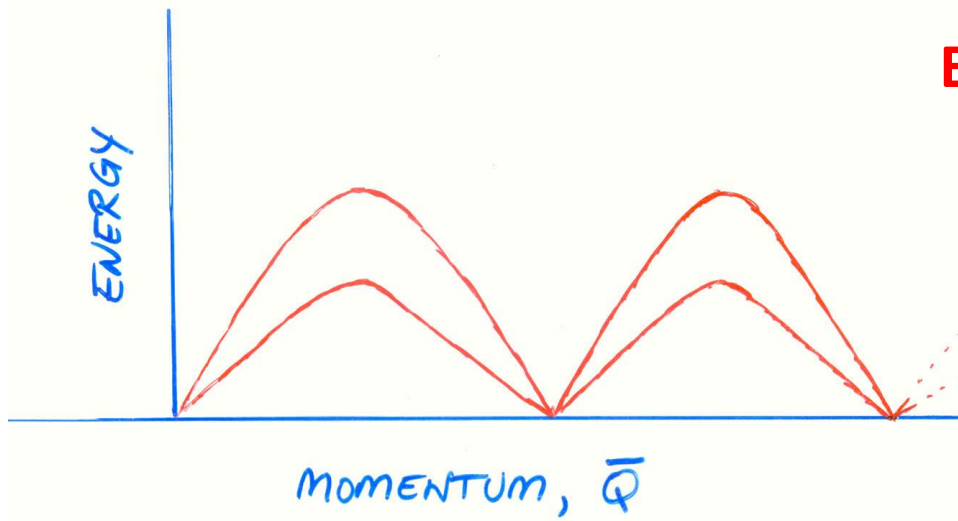
Remains fixed for any sample rotation



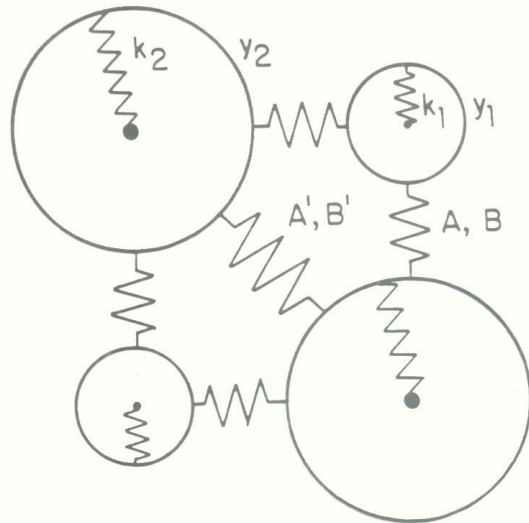




Elementary Excitations in Solids

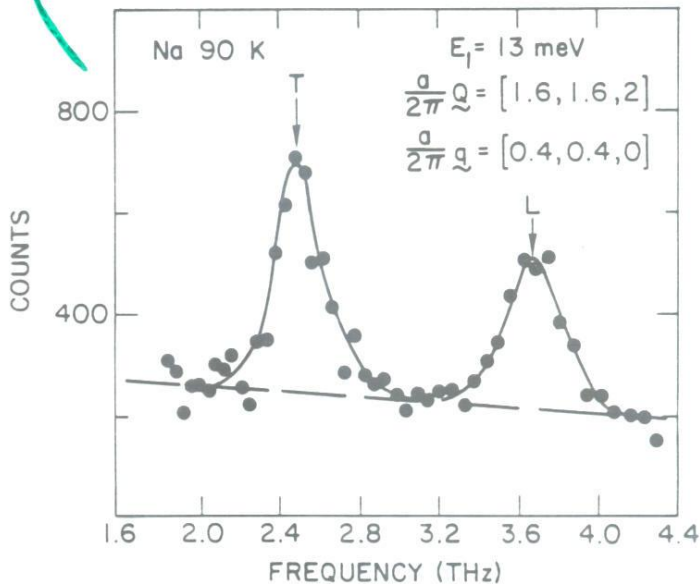
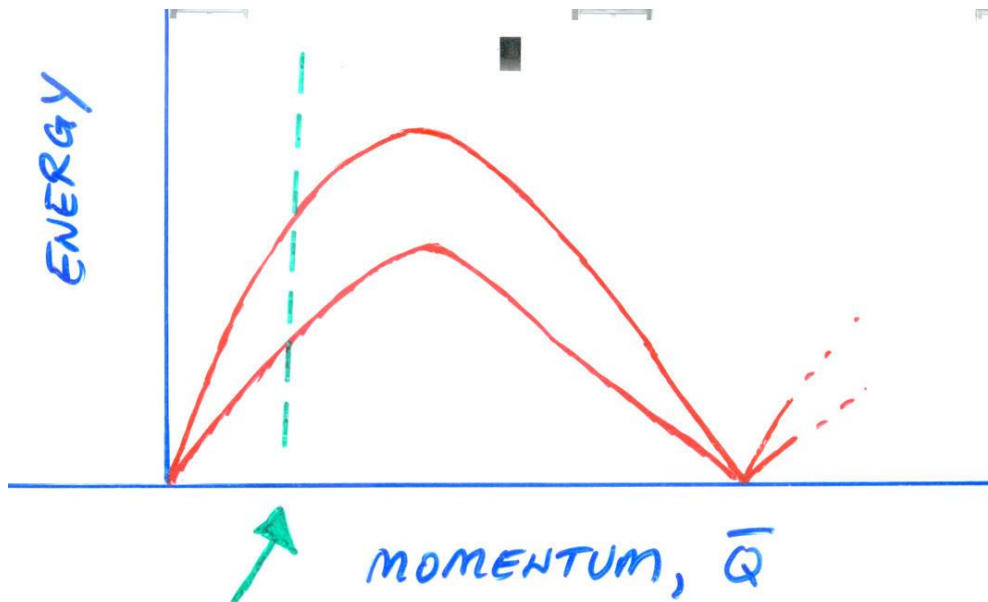


- Lattice Vibrations (Phonons)
- Spin Fluctuations (Magnons)



Energy vs Momentum

- Forces which bind atoms together in solids

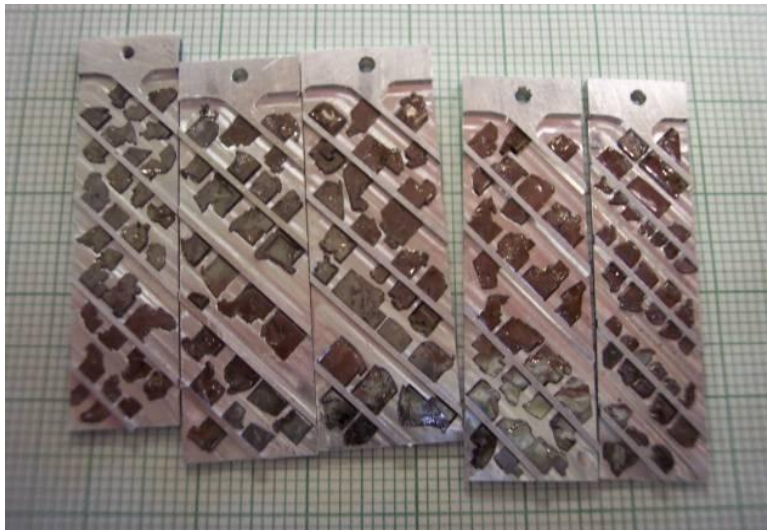


Constant Q, Constant E
3-axis technique allow us to
Put Q-Energy space on a grid,
And scan through as we wish

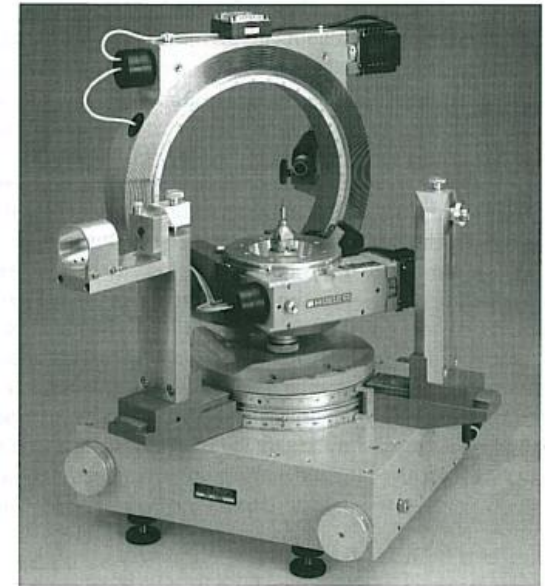
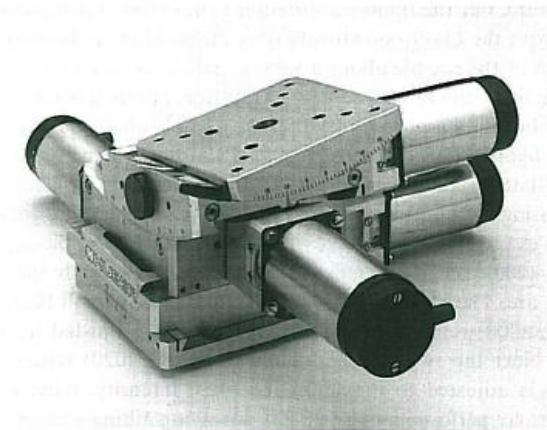
Map out elementary excitations
In Q-energy space (dispersion
Surface)

Samples

- **Samples need to be BIG**
 - ~ gram or cc
 - Counting times are long (mins/pt)
- **Sample rotation**
- **Sample tilt**

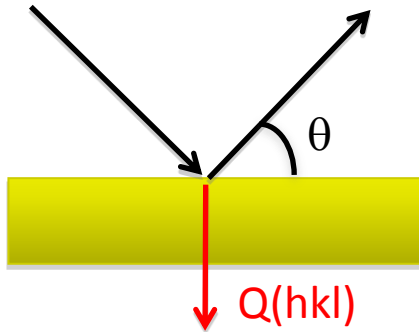


Co-aligned CaFe_2As_2 crystals

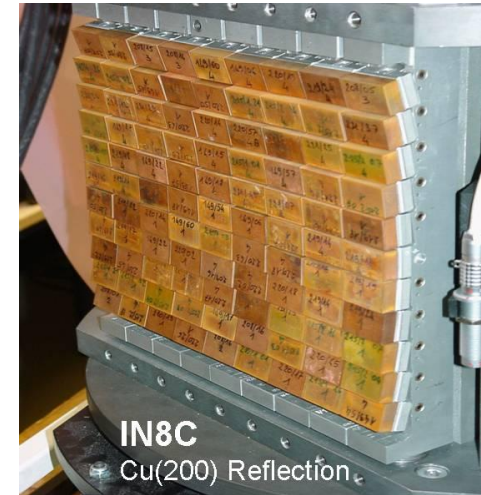


Monochromators

- Selects the incident wavevector



$$Q(hkl) = \frac{2\pi}{d(hkl)} = 2k_i \sin \theta$$



- Reflectivity
- focusing
- high-order contamination
eg. $\lambda/2$ PG(004)

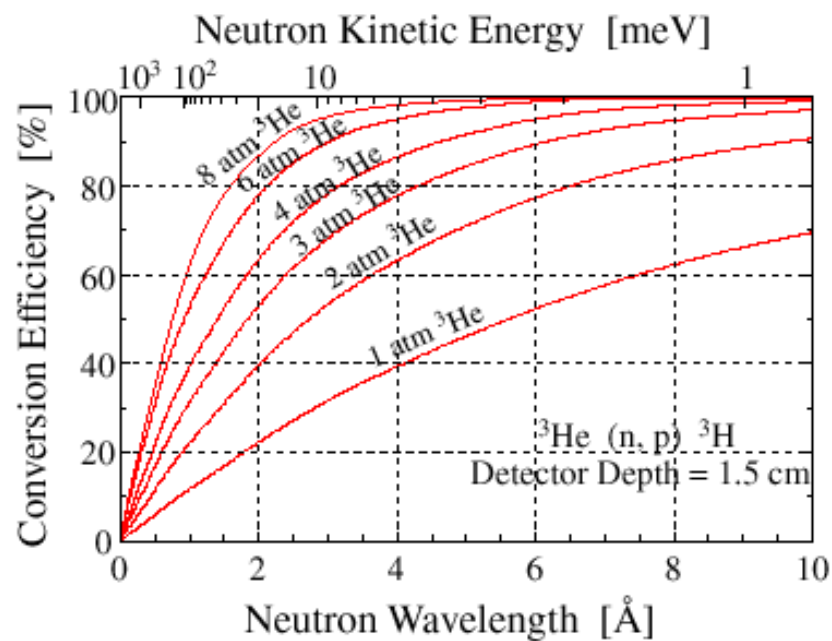
Mono	d(hkl)	uses
PG(002)	3.353	General
Be(002)	1.790	High k_i
Si(111)	3.135	No $\lambda/2$

Detectors



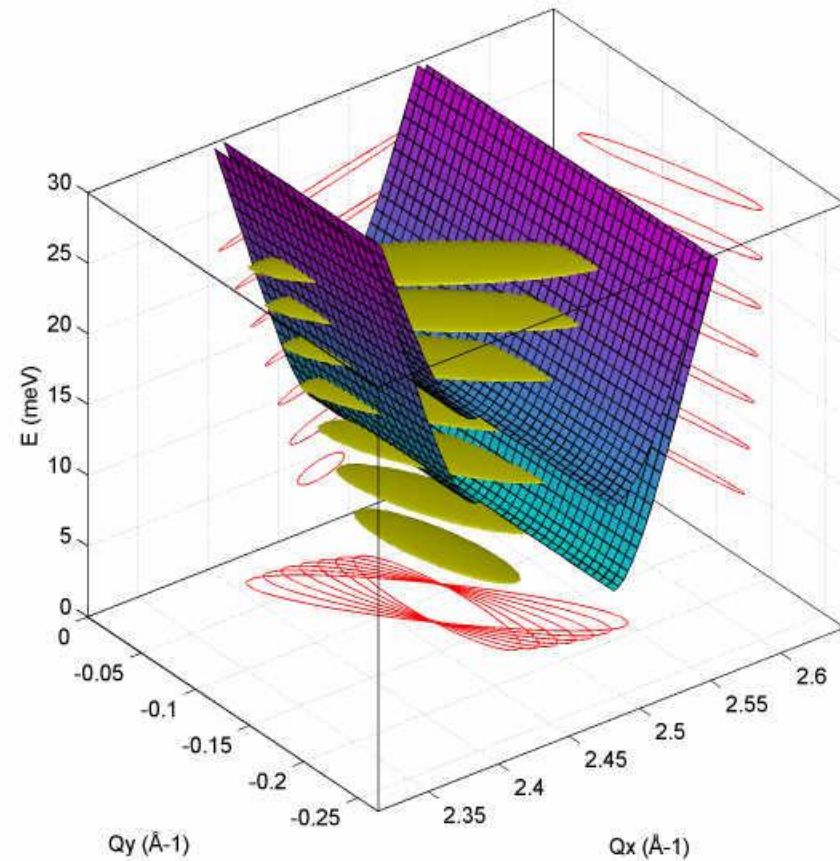
- **Gas Detectors**
- $n + {}^3\text{He} \rightarrow {}^3\text{H} + p + 0.764 \text{ MeV}$
- Ionization of gas
- e^- drift to high voltage anode
- High efficiency

- **Beam monitors**
- Low efficiency detectors for measuring beam flux



Resolution

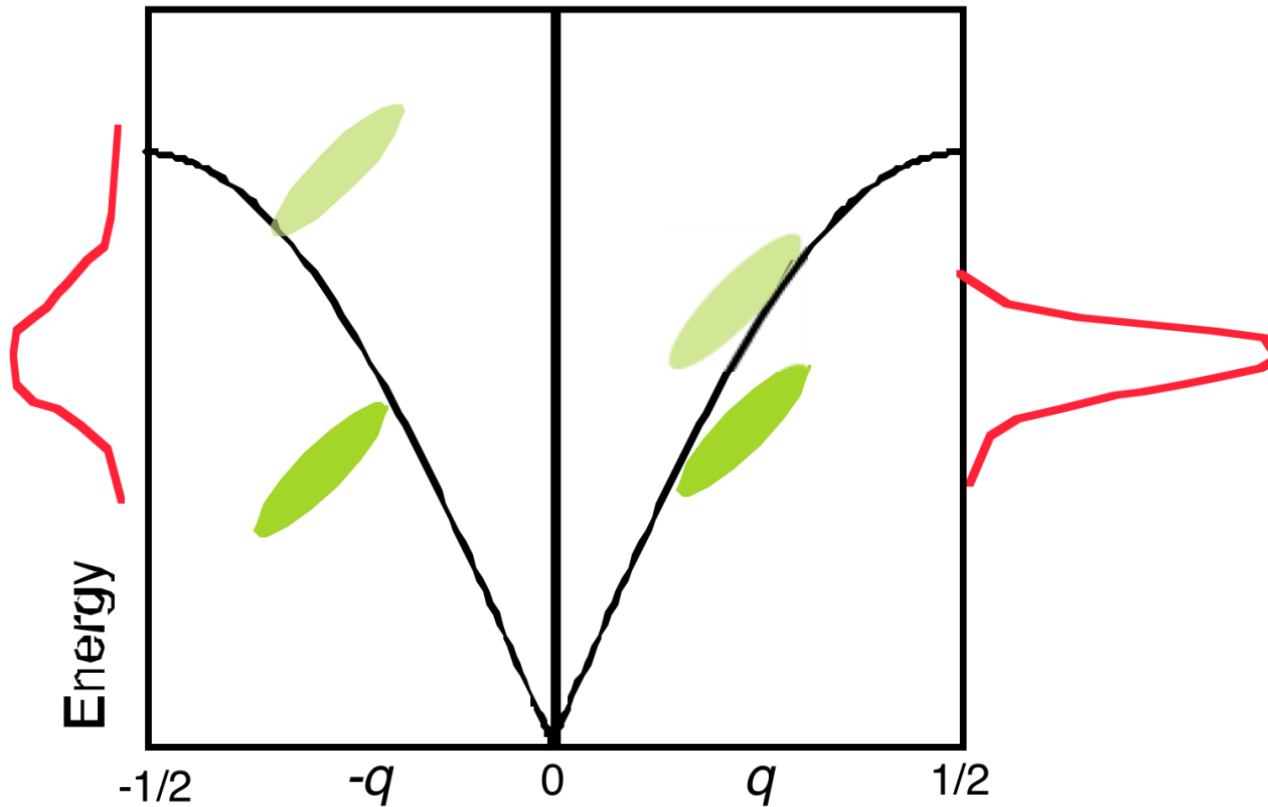
- **Resolution ellipsoid**
 - Beam divergences
 - Collimations/distances
 - Crystal mosaics/sizes/angles
- **Resolution convolutions**



$$I(\mathbf{Q}_0, \omega_0) = \int S(\mathbf{Q}_0, \omega_0) R(\mathbf{Q} - \mathbf{Q}_0, \omega - \omega_0) d\mathbf{Q} d\omega$$

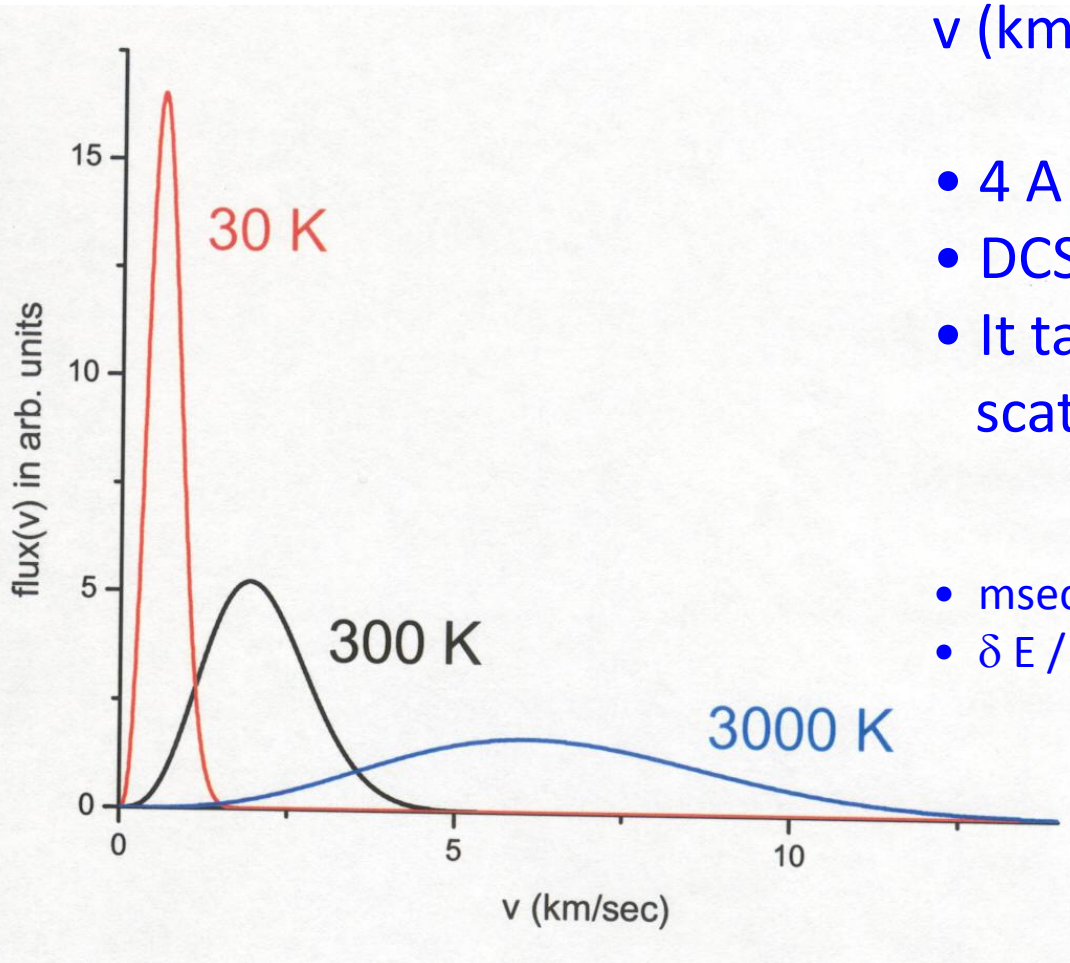
Resolution focusing

- Optimizing peak intensity
- Match slope of resolution to dispersion



Neutrons have *mass*

so higher energy means faster – lower energy means slower



$$v \text{ (km/sec)} = 3.96 / \lambda \text{ (Å)}$$

- 4 Å neutrons move at ~ 1 km/sec
- DCS: 4 m from sample to detector
- It takes 4 msec for elastically scattered 4 Å neutrons to travel 4 m
- msec timing of neutrons is easy
- $\delta E / E \sim 1\text{-}3\%$ - very good !

We can measure a neutron's energy, wavelength by measuring its *speed*

Time-of-flight methods



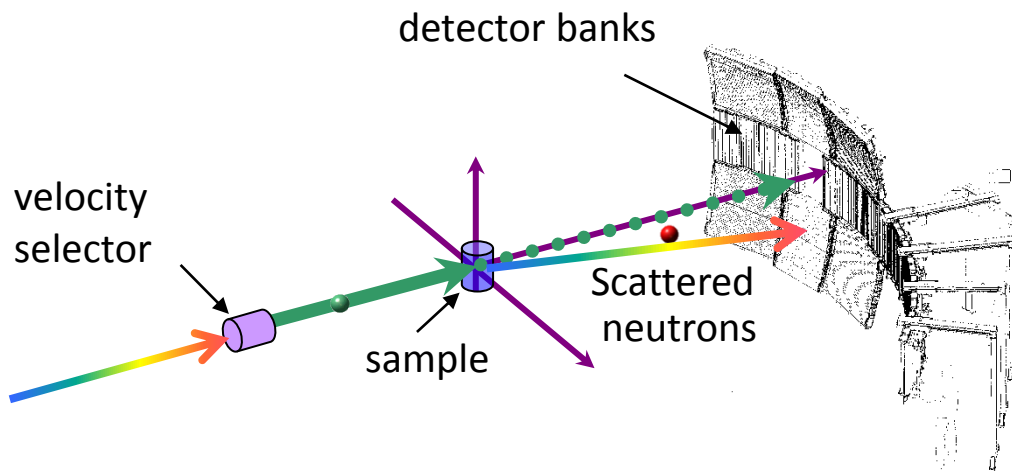
Spallation neutron source



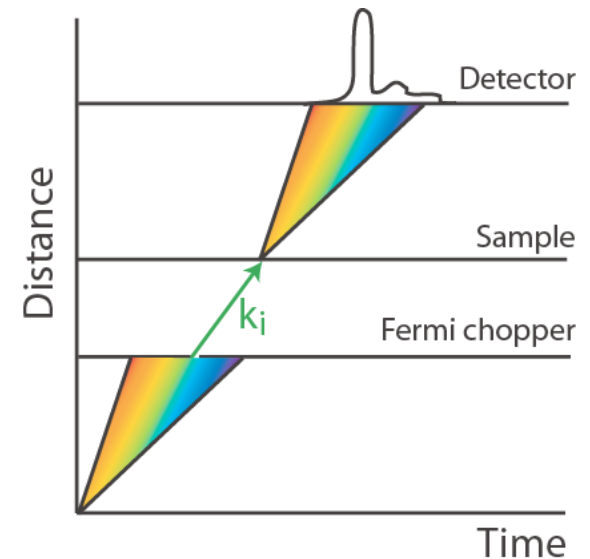
Pharos – Lujan Center

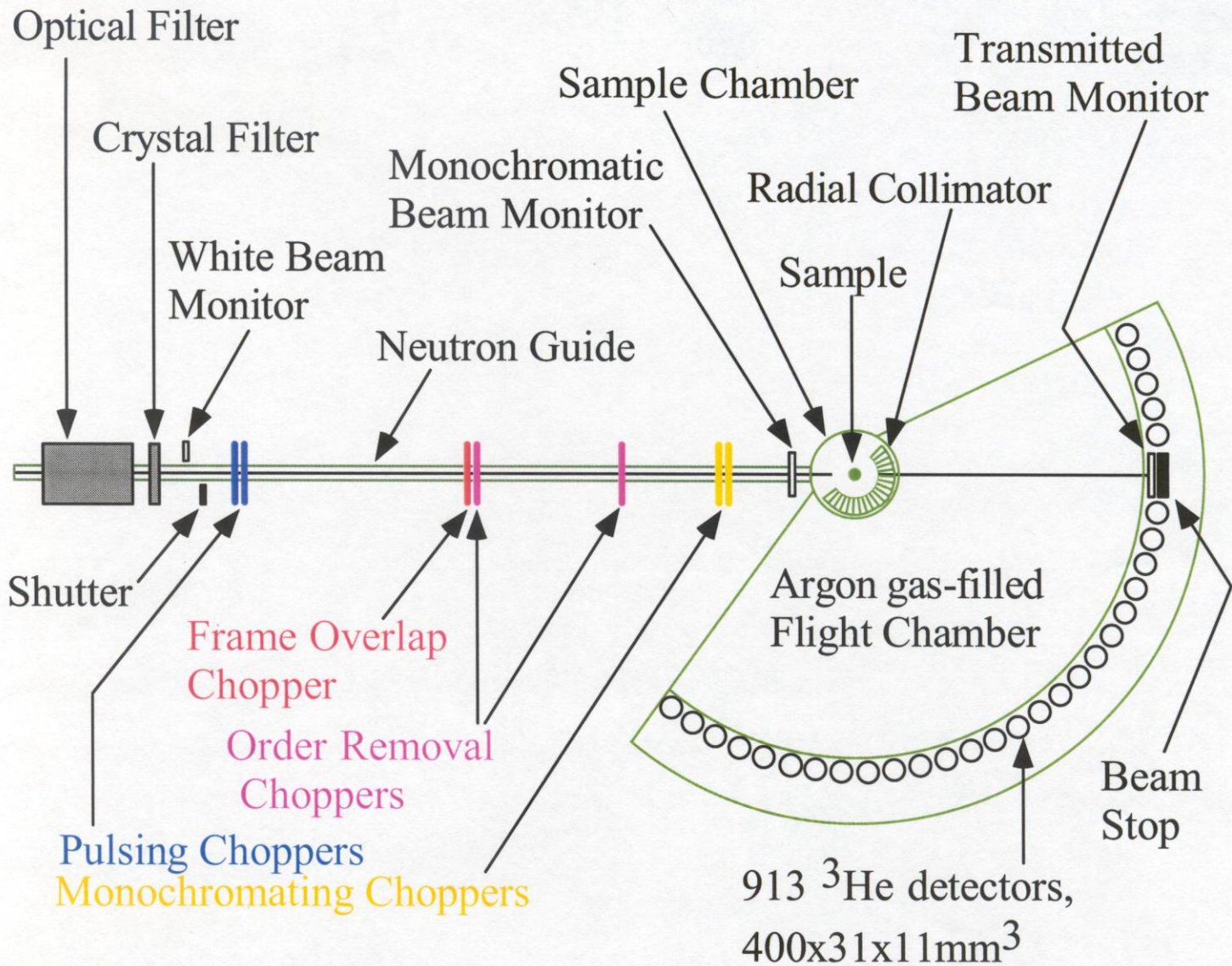
- Effectively utilizes time structure of pulsed neutron groups

$$t = \frac{d}{v} = \left(\frac{m}{h} d \right) \lambda$$

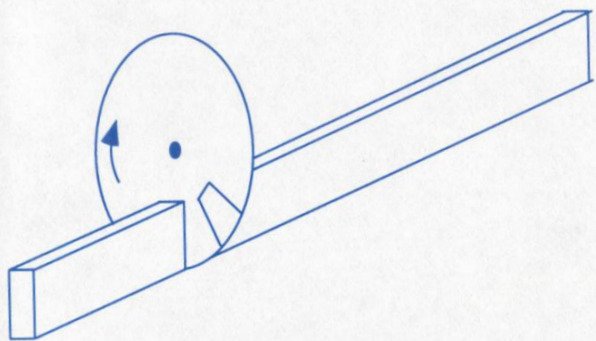


NXS School

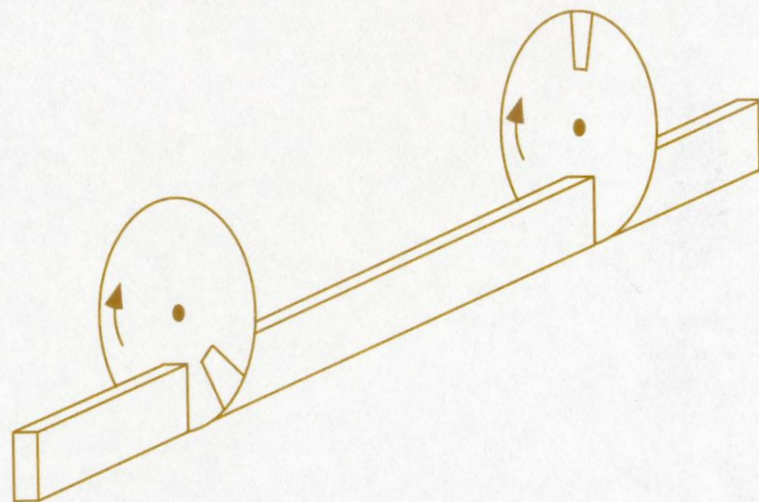




A single (disk) chopper pulses the neutron beam.



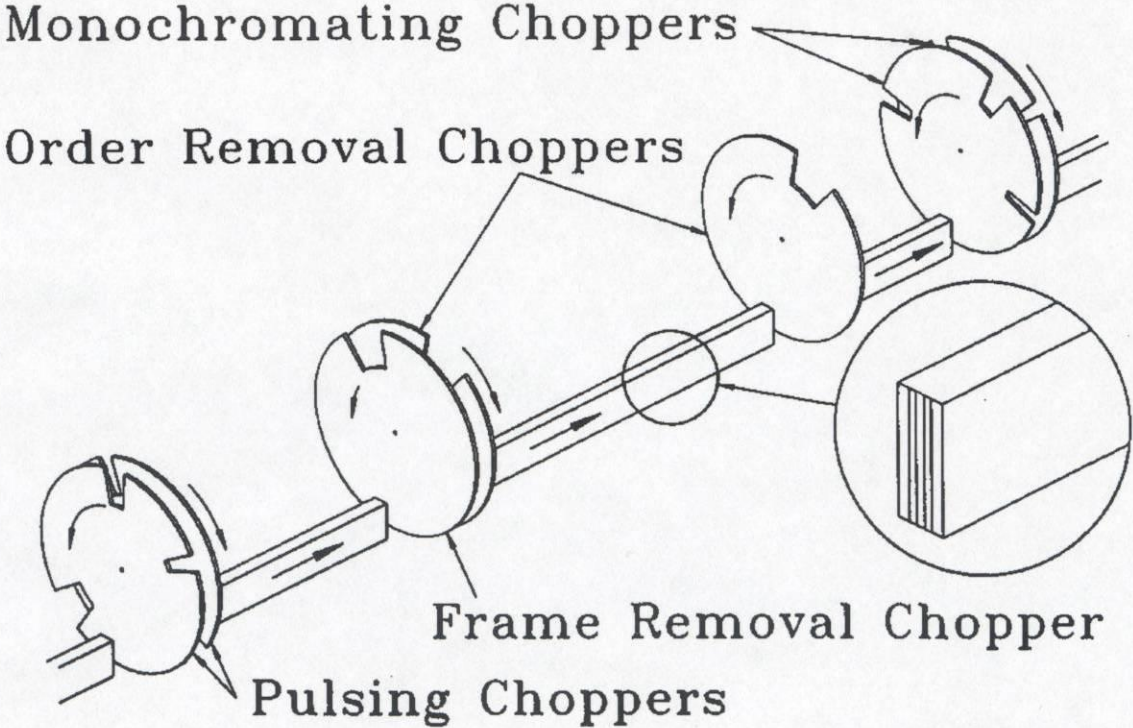
A second chopper selects neutrons within a narrow range of speeds.



Counter-rotating choppers (close together), with speed \diamond , behave like single choppers with speed $2\diamond$. They can also permit a choice of pulse widths.

Additional choppers remove “contaminant” wavelengths and reduce the pulse frequency at the sample position.

The DCS has seven choppers, 4 of which have 3 “slots”



Disk 4B



Fermi Choppers

- Body radius ~ 5 cm
- Curved absorbing slats
 - B or Gd coated
 - \sim mm slit size
- $f = 600$ Hz (max)
- Acts like shutter, $\Delta t \sim \mu\text{s}$

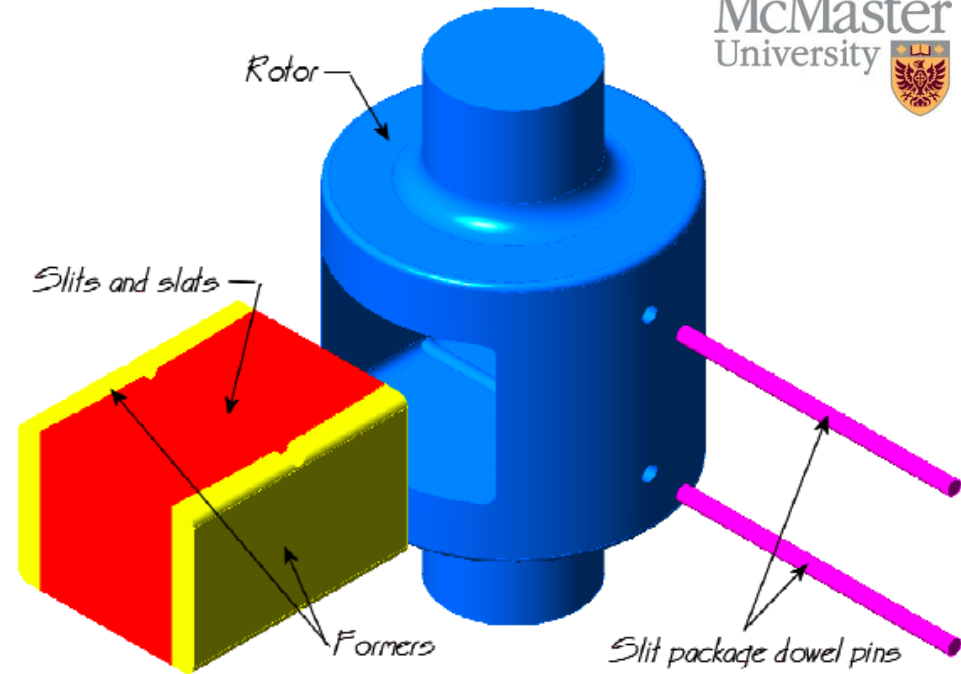
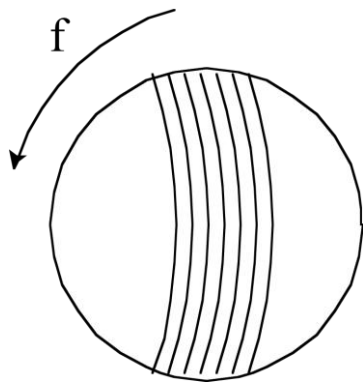
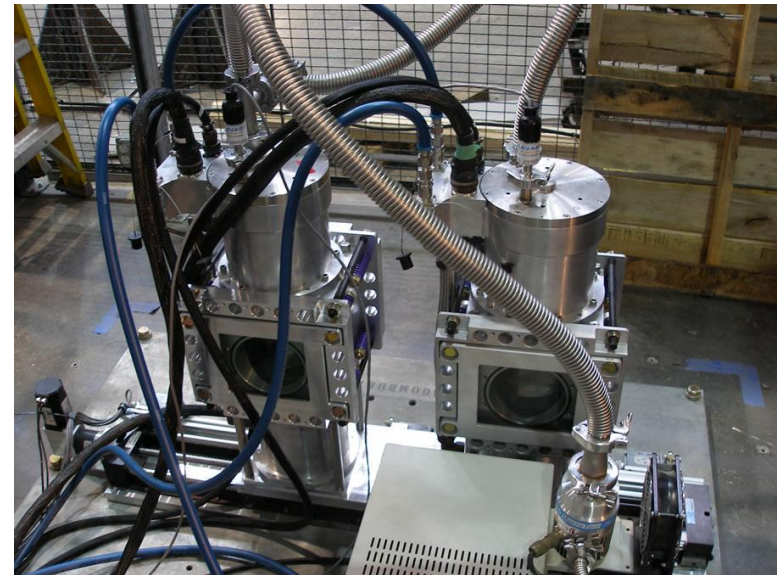
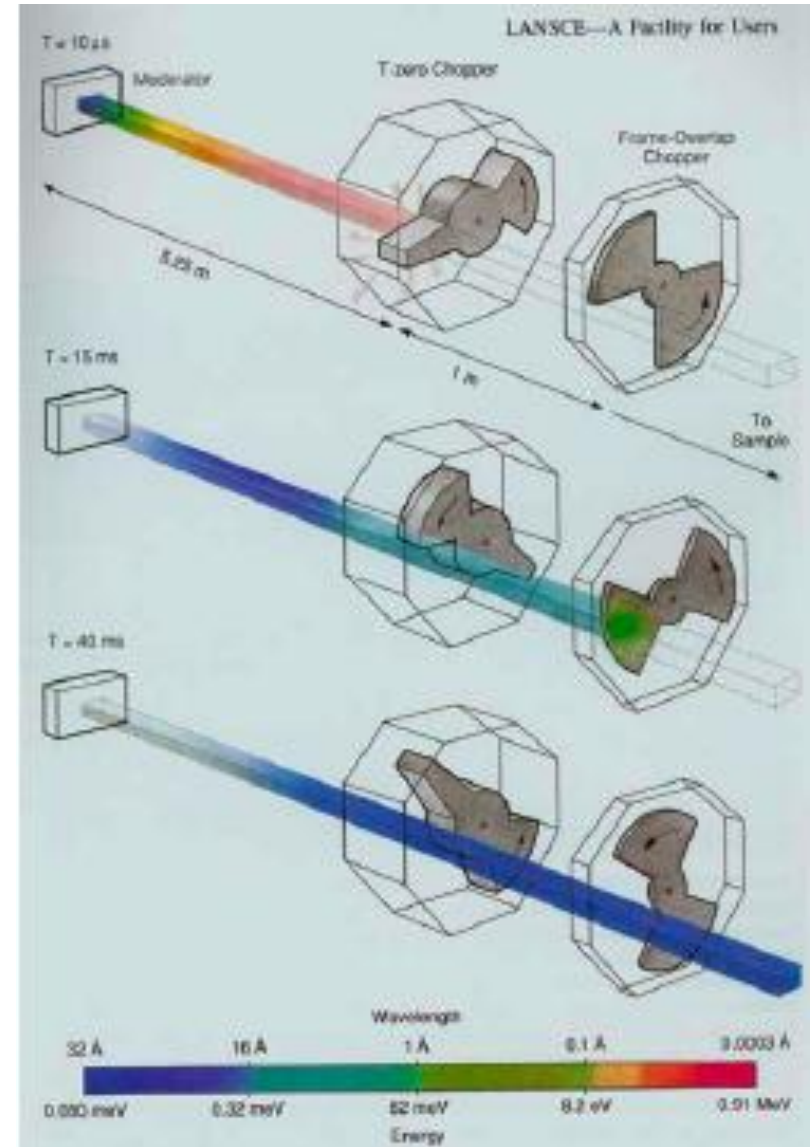
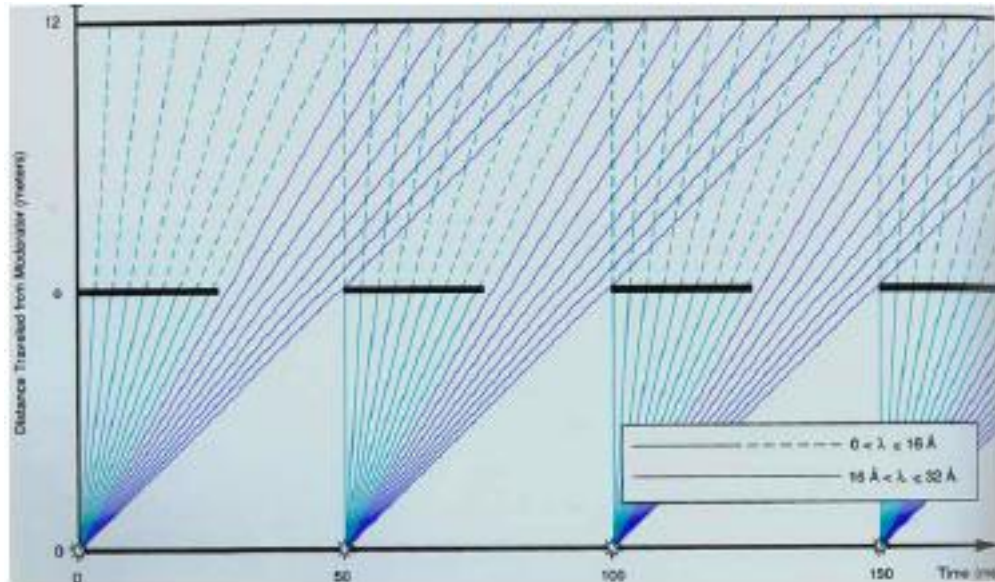


Figure 1. ISIS MAPS chopper and slit package assembly – exploded view



T-zero chopper

- Background suppression
- Blocks fast neutron flash

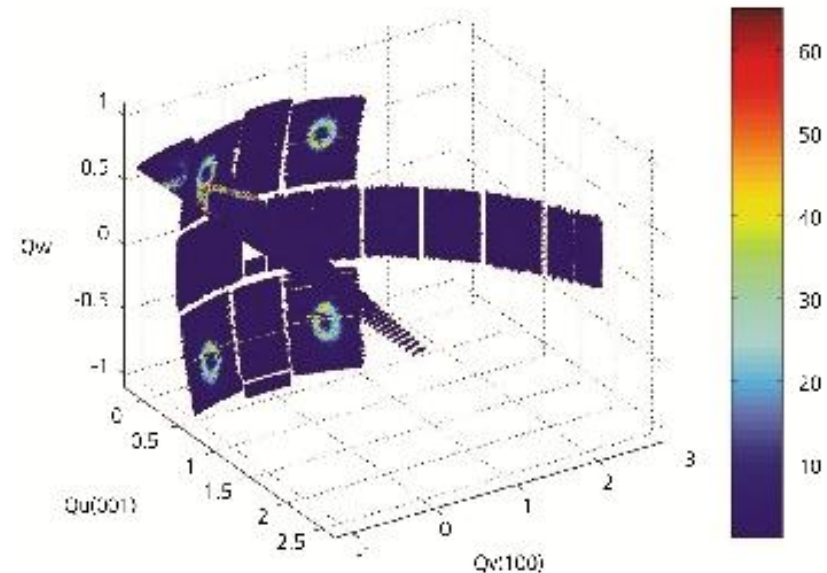


Position sensitive detectors

- ^3He tubes (usu. 1 meter)
- Charge division
- Position resolution \sim cm
- Time resolution \sim 10 ns



MAPS detector bank



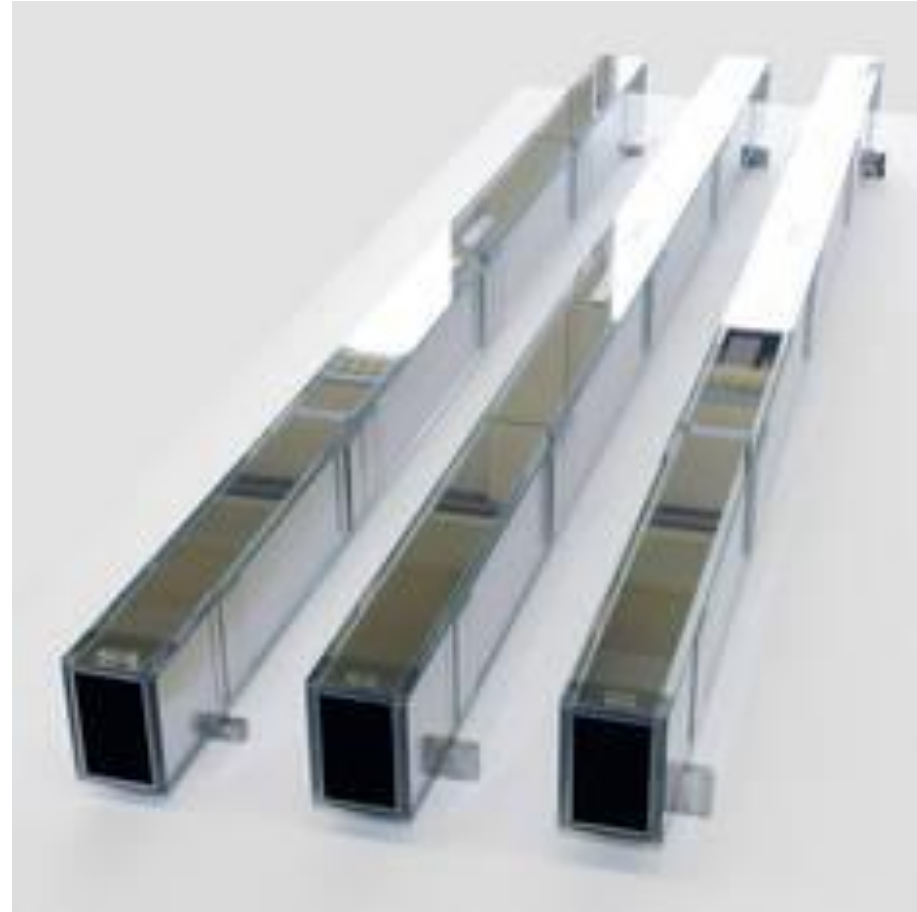
Sample environment

- Temperature, field, pressure
- Heavy duty for large sample environment
 - CCR
 - He cryostats
 - SC magnets
 - ...
- Can be machined from Al
 - ~ neutron transparent
 - relatively easy to work with

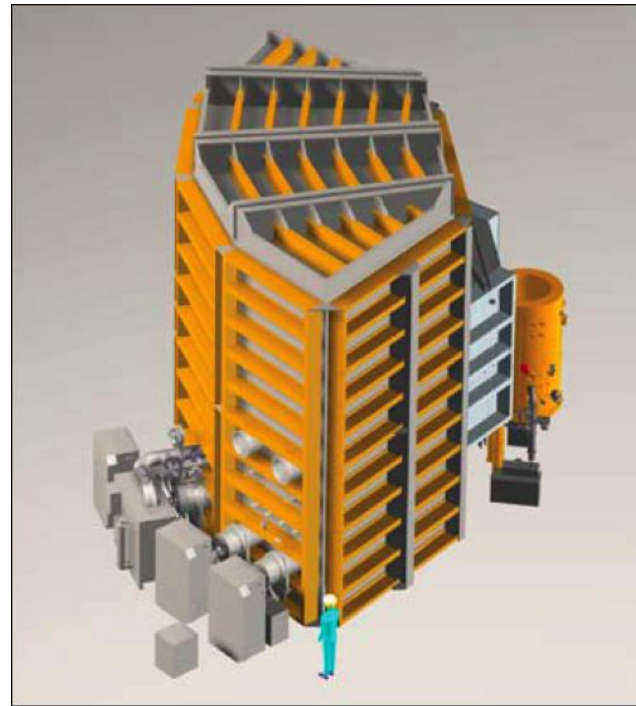


Guides

- Transport beam over long distances
- Background reduction
- Total external reflection
 - Ni coated glass
 - Ni/Ti multilayers (supermirror)



Size matters



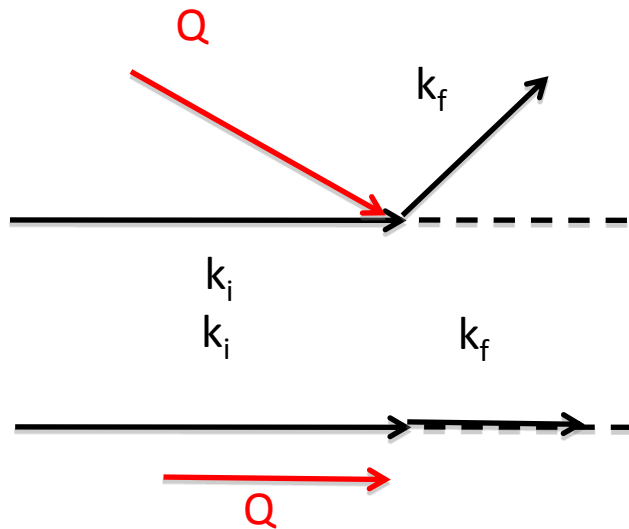
SEQUOIA detector
vacuum vessel

- **Length = resolution**
 - Instruments $\sim 20 - 40$ m long
 - E-resolution $\sim 2-4\% E_i$
- **More detectors**
 - SEQUOIA – 1600 tubes, 144000 pixels
 - Solid angle coverage 1.6 steradians
- **Huge data sets**
- 0.1 – 1 GB

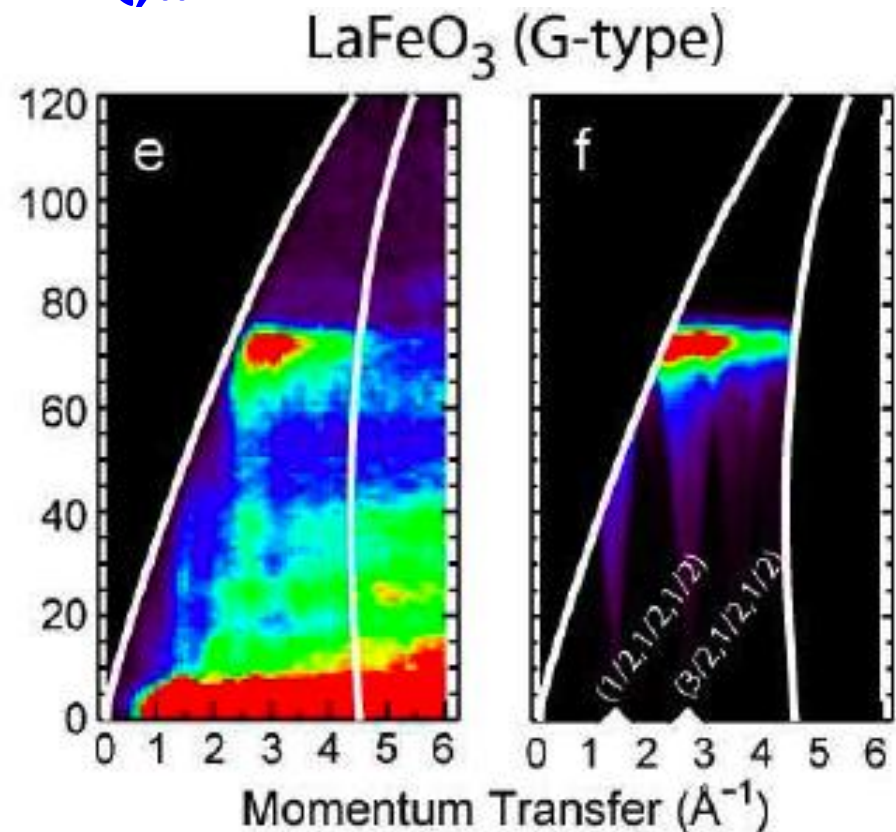


Kinematic limitations

- Many combinations of k_i, k_f for same Q, ω
 - Only certain configurations are used (eg. E_f -fixed)
- Cannot “close triangle” for certain Q, ω due to kinematics

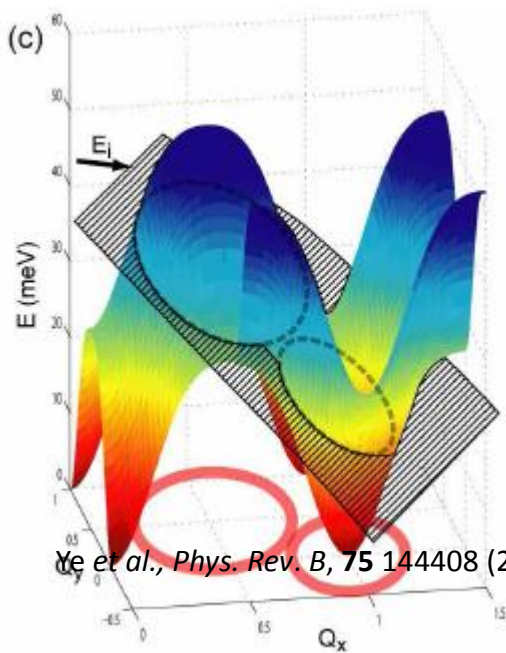


Minimum accessible Q

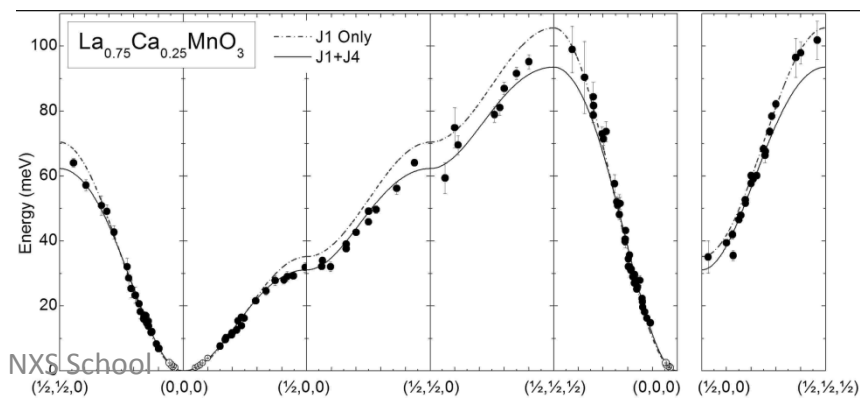
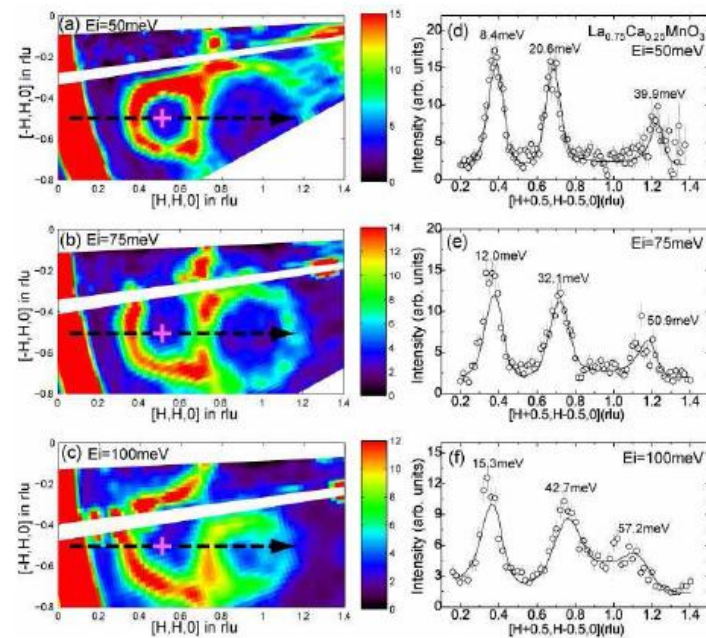


Data visualization

- Large, complex data from spallation sources
- Measure $S(\mathbf{Q}, \omega)$ – 4D function



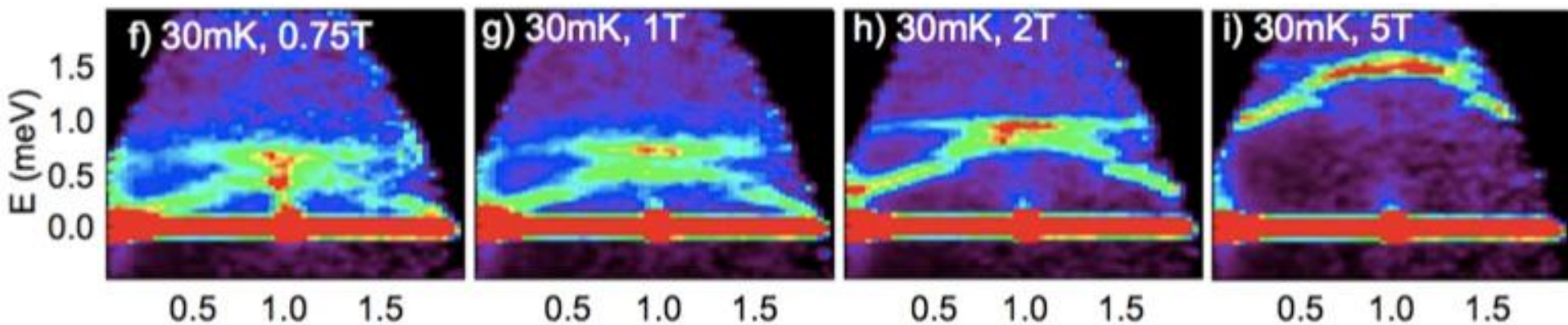
Ye et al., *Phys. Rev. B*, **75** 144408 (2007).



Field-induced order in the Pyrochlore $\text{Yb}_2\text{Ti}_2\text{O}_7$:

Weak magnetic field // $[110]$ induces LRO

*appearance of long-lived spin waves
at low T and moderate H*



References

General neutron scattering

G. Squires, “Intro to theory of thermal neutron scattering”, Dover, 1978.

S. Lovesey, “Theory of neutron scattering from condensed matter”, Oxford, 1984.

R. Pynn, <http://www.mrl.ucsb.edu/~pynn/>.

Polarized neutron scattering

Moon, Koehler, Riste, Phys. Rev **181**, 920 (1969).

Triple-axis techniques

Shirane, Shapiro, Tranquada, “Neutron scattering with a triple-axis spectrometer”, Cambridge, 2002.

Time-of-flight techniques

B. Fultz, http://www.cacr.caltech.edu/projects/danse/ARCS_Book_16x.pdf